

# **Anharmonic lattice dynamics of clathrates explained by vibrational dynamical mean-field theory**

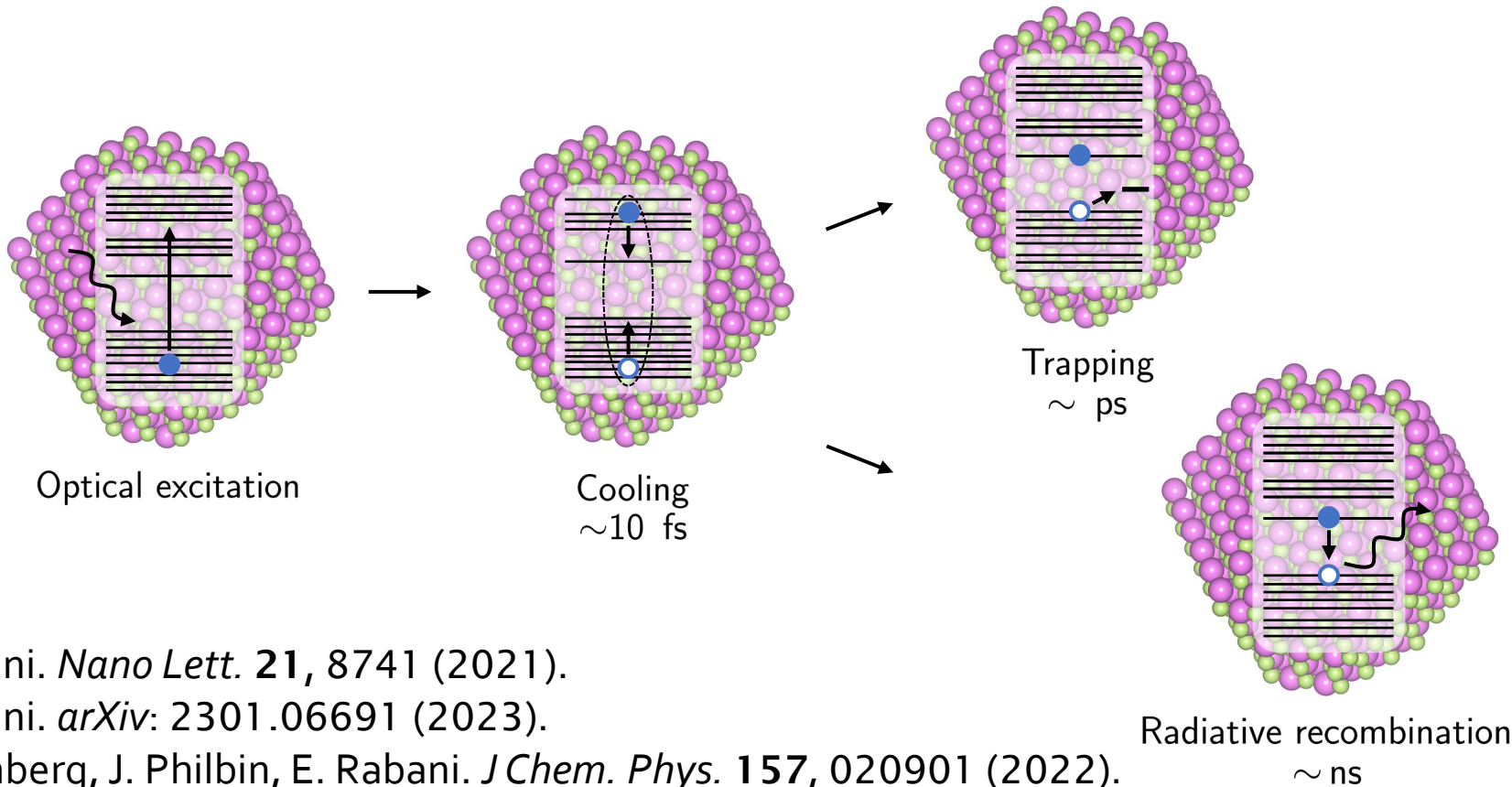
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**Berkelbach Group, Columbia University**

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# PhD: How are optoelectronic properties of QDs affected by atomic vibrations?

- Understanding exciton-phonon coupling and its impact on processes such as exciton cooling, photoluminescence, and trapping



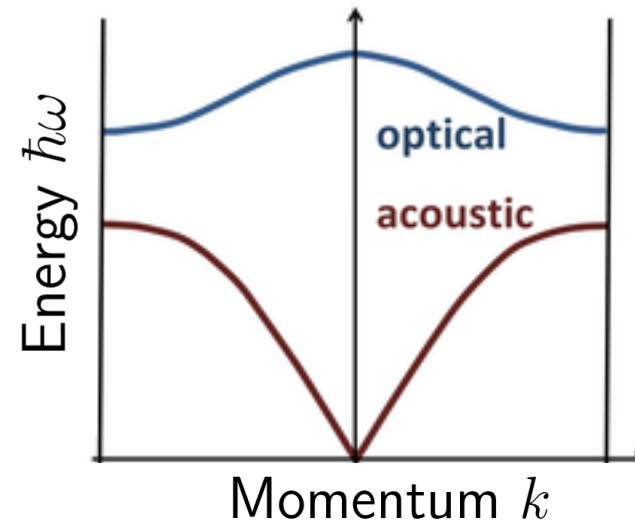
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D. Jasrasaria, E. Rabani. *arXiv*: 2301.06691 (2023).

D. Jasrasaria, D. Weinberg, J. Philbin, E. Rabani. *J Chem. Phys.* **157**, 020901 (2022).

# Postdoc: How can we rationally control vibrational properties of materials?

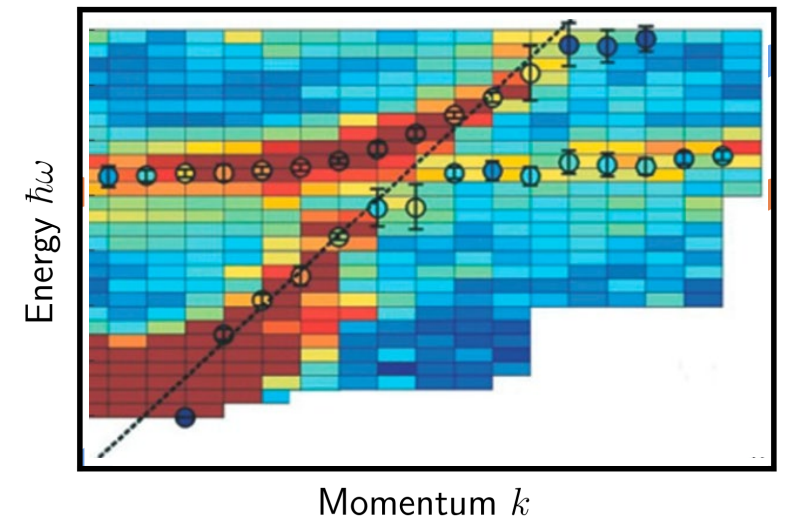
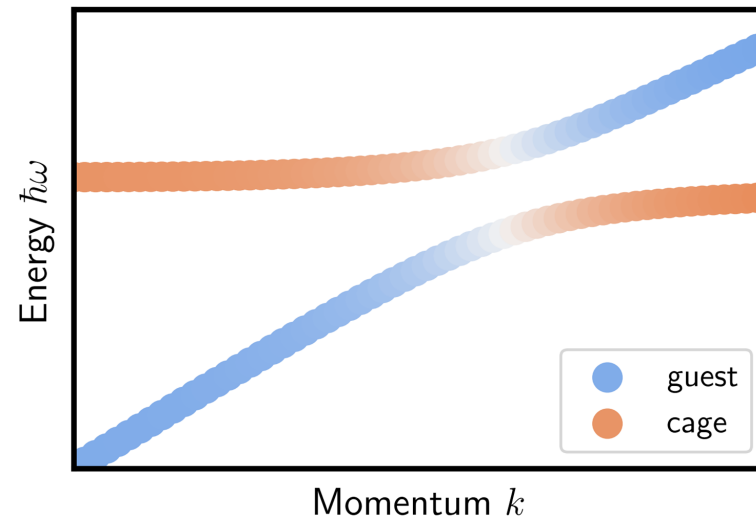
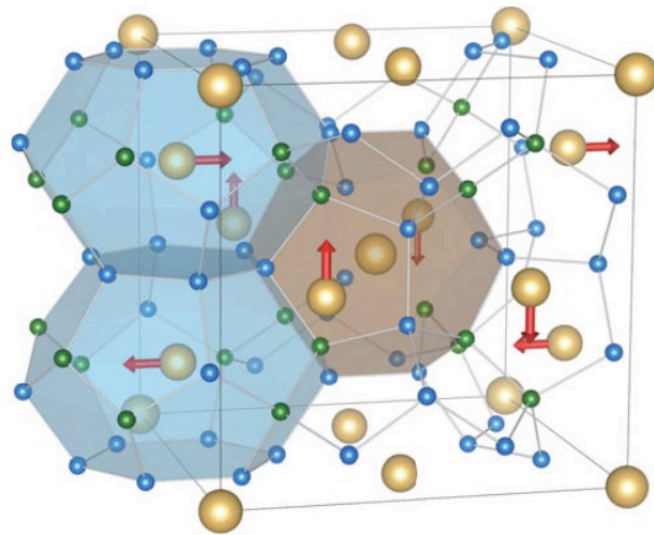
- Simplest picture: collection of non-interacting phonons
  - Phonons are quantized, collective oscillations of the lattice
  - Carry sound and heat, and can interact with light



- Some materials have significant "anharmonicity" (phonon-phonon interactions)

# Thermoelectrics convert heat to electricity

- Good thermoelectrics have low thermal conductivities
  - Anharmonic vibrational structure (i.e., phonon-phonon interactions)
- Clathrates: cage-like structures with loosely bound, embedded guest
- Hybridization of cage acoustic mode and guest optical mode

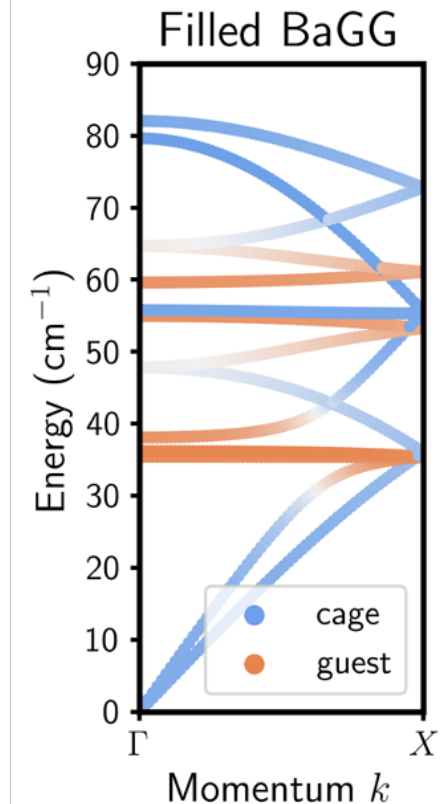
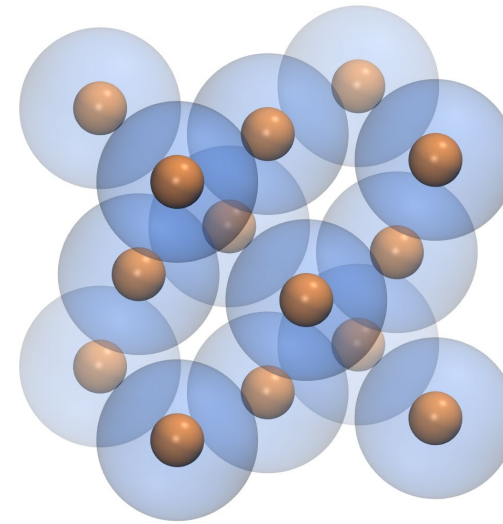
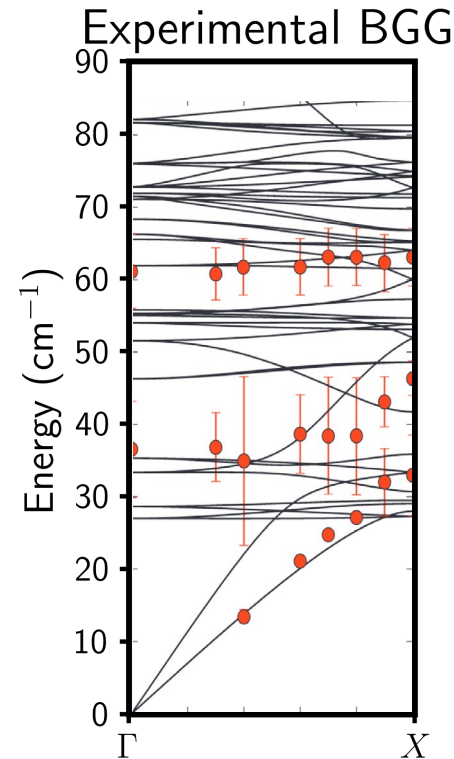
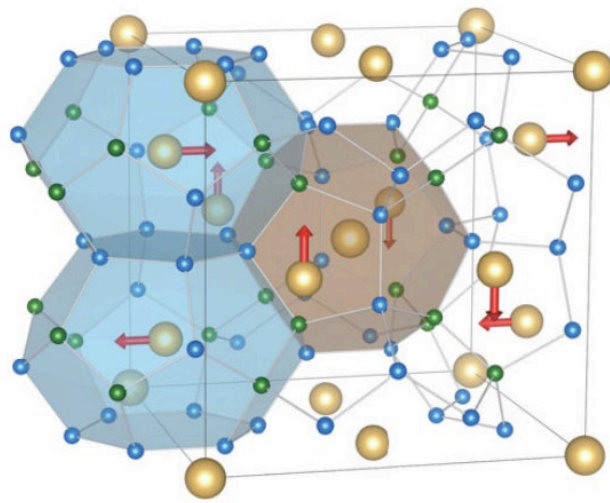


T. Tadano, S. Tsuneyuki. *Phys. Rev. Lett.* **120**, 105901 (2018).

M. Christensen et al. *Nature Mater.* **7**, 811 (2008).

# Developing a coarse-grained clathrate model: $\text{Ba}_8\text{Ga}_{16}\text{Ge}_{30}$ and $\text{Sr}_8\text{Ga}_{16}\text{Ge}_{30}$

- Lattice of large “cage” atoms with Lennard-Jones interactions
- Small “guest” atom with quartic interaction with its own cage



J. Dong et al. *J. Appl. Phys.* **87**, 7726 (2000). Momentum  $k$

T. Tadano et al. *Phys. Rev. Lett.* **114**, 095501 (2015).

# Accounting for anharmonicity

- Anharmonic single-particle phonon Green's function:

$$i\hbar D_{\lambda\lambda'}(\mathbf{k}, \omega) = \int_0^\infty dt e^{i\omega t} \langle [u_\lambda(\mathbf{k}, t), u_{\lambda'}(-\mathbf{k}, 0)] \rangle$$

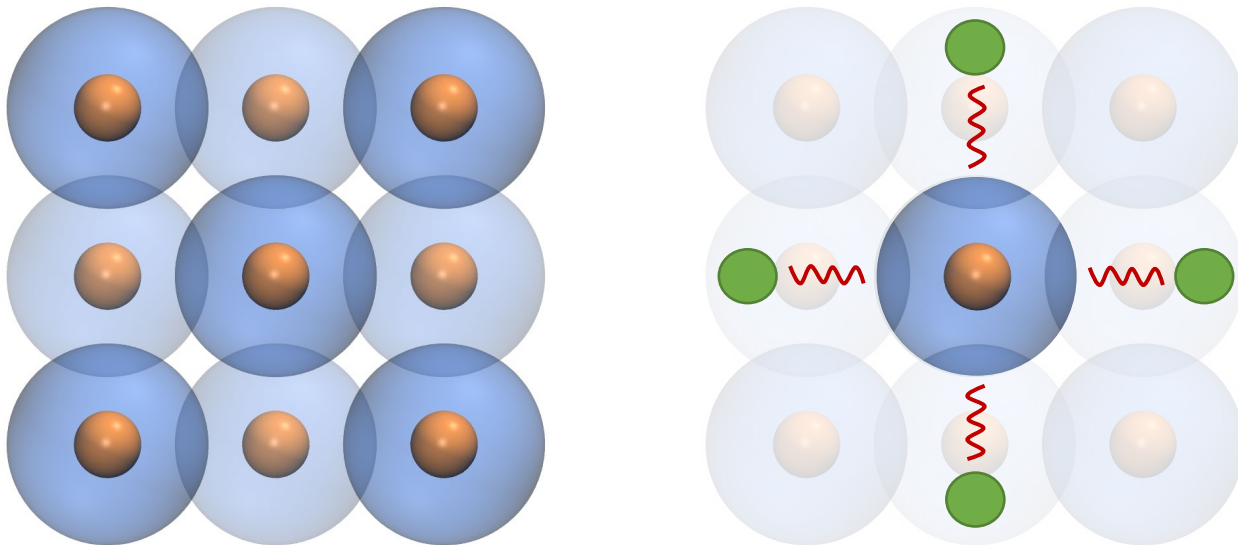
- Could calculate using MD... challenging!

- Write anharmonic GF in terms of noninteracting GF and self-energy

$$\mathbf{D}^{-1}(\mathbf{k}, \omega) = \mathbf{D}_0^{-1}(\mathbf{k}, \omega) - \boldsymbol{\pi}(\mathbf{k}, \omega)$$

# VDMFT: Map lattice dynamics onto impurity problem

- Nonperturbatively solve for self-energy of a finite **system coupled to a harmonic bath**
  - Key approximation of local self-energy:  $\pi(\mathbf{k}, \omega) \approx \pi(\omega)$



$$H_s = \sum_{\alpha} \frac{p_{\alpha}^2}{2m_{\alpha}} + V_{\text{loc}}(\{\mathbf{R}_{\alpha}\})$$

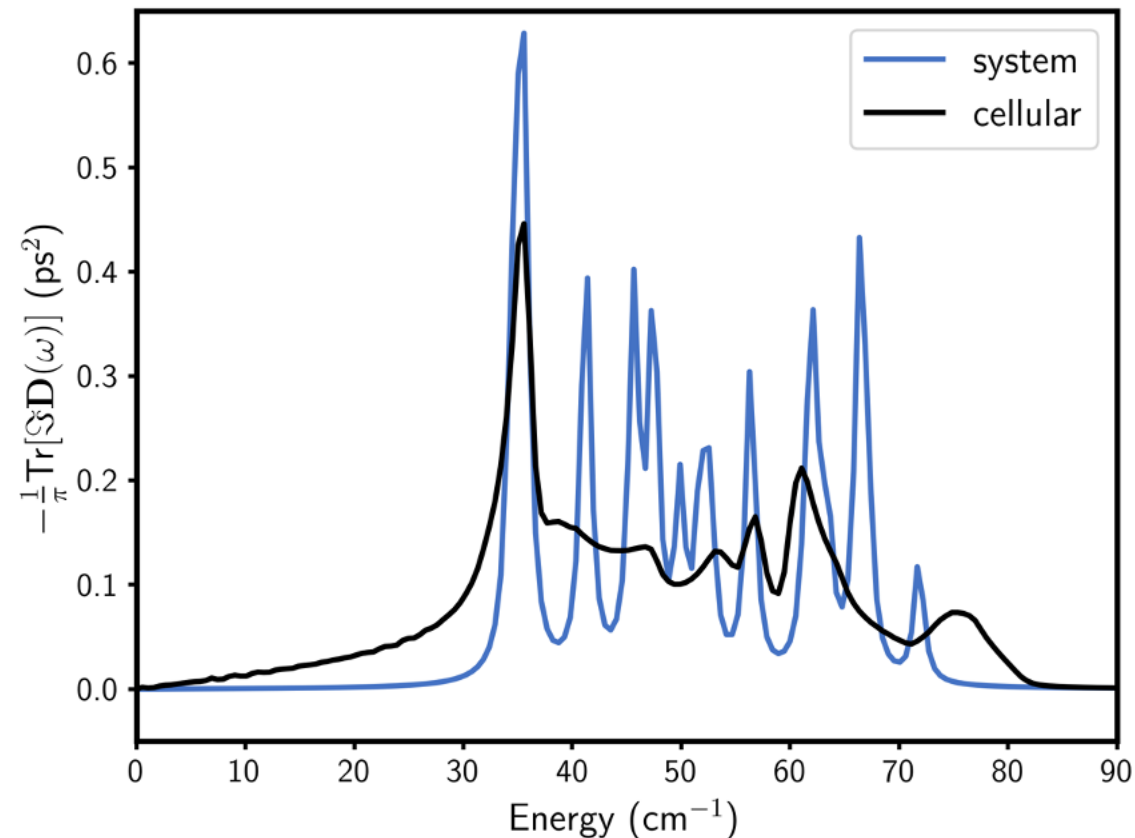
$$H_b = \frac{1}{2} \sum_m (p_m^2 + \omega_m^2 x_m^2)$$

$$H_{sb} = \sum_{\alpha i, m} c_{\alpha i, m} u_{\alpha i} x_m$$

# 1) Defining the impurity problem

- Difference between the **lattice GF** and **system GF** defines **system-bath coupling**:

$$\Delta(\omega) = D_{\text{sys}}^{-1}(\omega) - D_C^{-1}(\omega)$$





## 2) Calculating impurity GF

- Impurity GF (anharmonic system + harmonic bath):

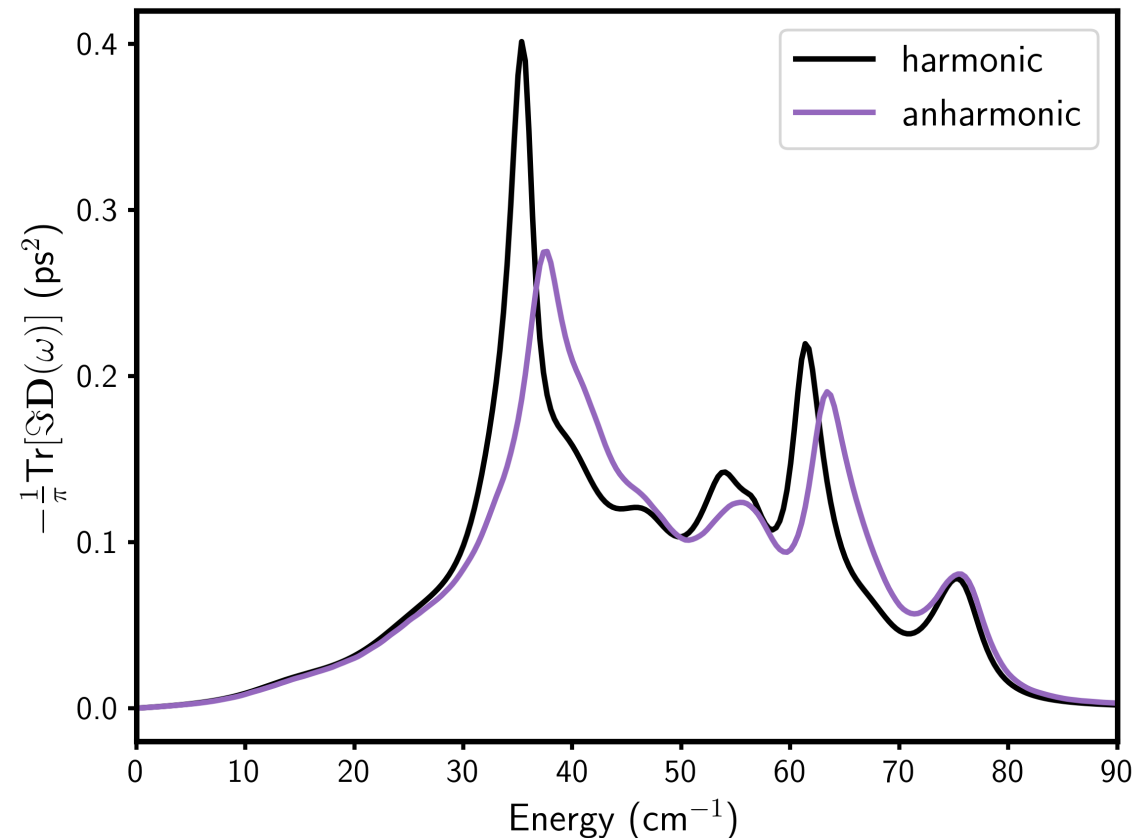
$$i\hbar D_{\text{imp}}(\omega) = \int_0^\infty dt e^{i\omega t} \langle [u(t), u(0)] \rangle$$

- Use generalized Langevin equation to solve system dynamics
  - Non-Markovian (frequency-dependent) **friction** makes it **dynamical** MFT

$$\ddot{u}(t) = -\frac{dV}{du} - \int_0^t ds \gamma(t-s) \dot{u}(s) + \xi(t)$$

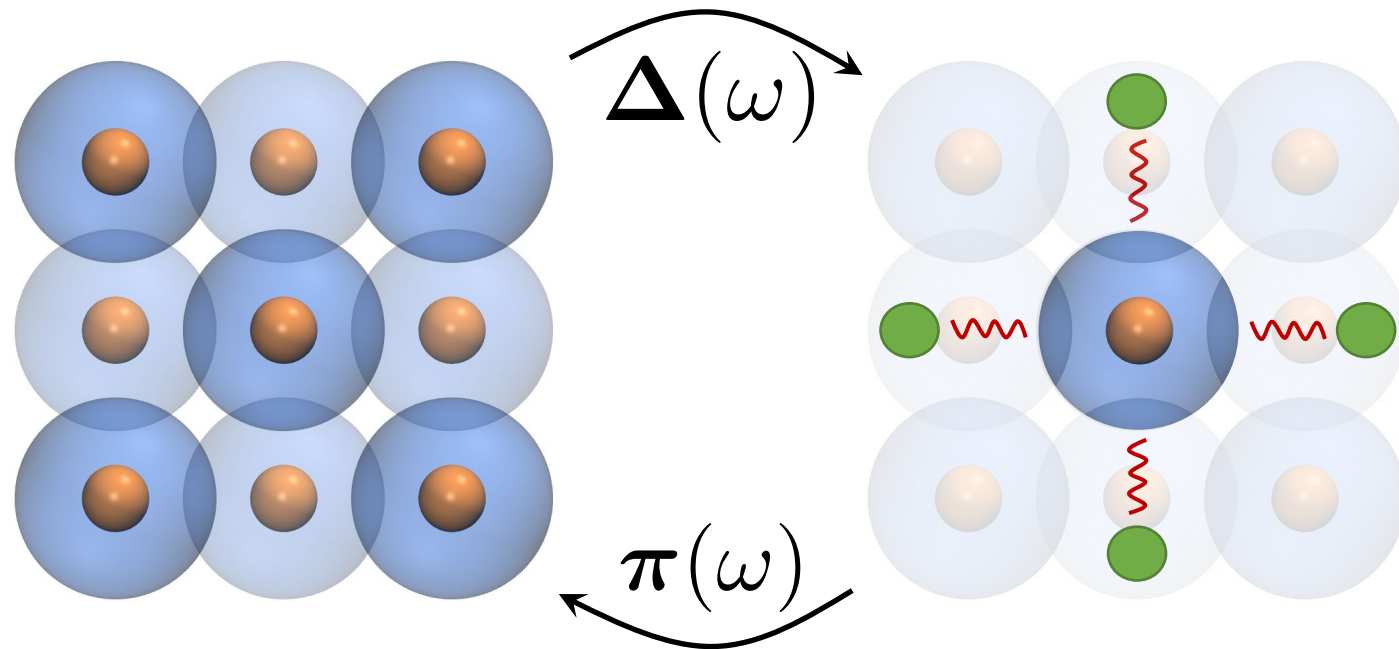
### 3) Getting (local) self-energy

- Difference between impurity **anharmonic GF** and **harmonic GF** is the **self-energy**:  
$$\pi(\omega) = D_{\text{imp},0}^{-1}(\omega) - D_{\text{imp}}^{-1}(\omega)$$



## 4) Iterate (if needed...)

- Use self-energy to define lattice GF that includes anharmonicity
  - New cellular GF, new **system** GF, new **hybridization**
  - New impurity GF and GF0, new self energy
- Eventually, convergence when  $D_C(\omega) = D_{\text{imp}}(\omega)$



# Anharmonic spectral function

$$D^{-1}(\mathbf{k}, \omega) = D_0^{-1}(\mathbf{k}, \omega) - \pi(\mathbf{k}, \omega)$$

# Anharmonic spectral function

$$A(\mathbf{k}, \omega) = -\frac{1}{\pi} \text{Tr}[\Im \mathbf{D}(\mathbf{k}, \omega)]$$

isolated phonon

finite lifetimes

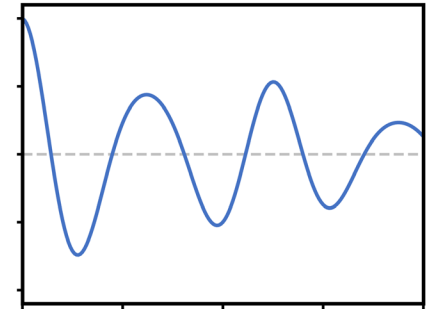
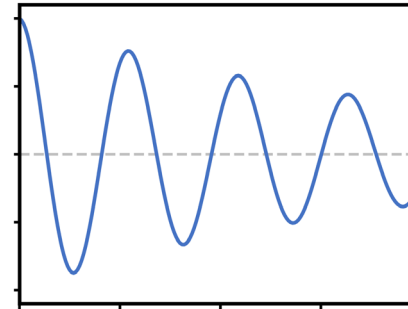
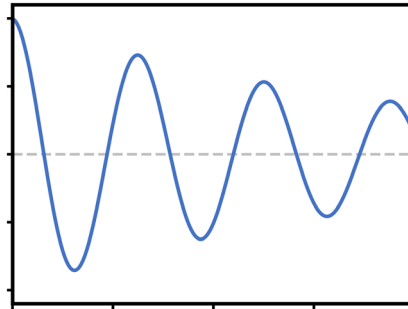
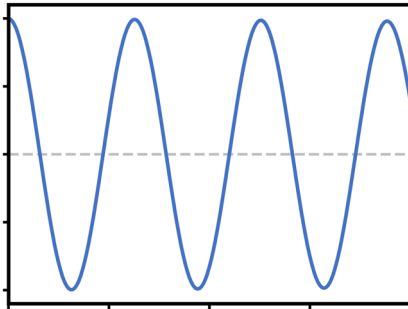
frequency shift

complicated...!

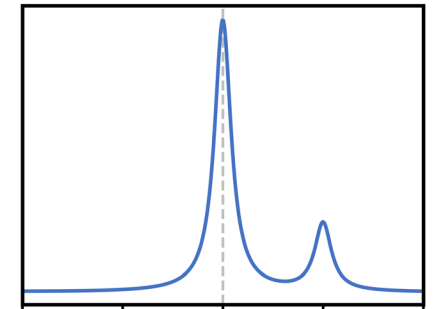
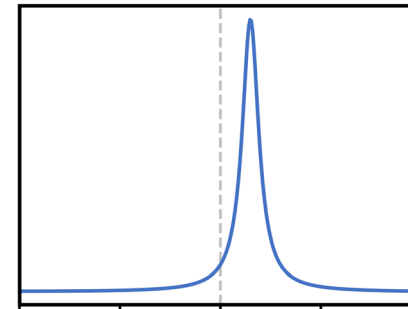
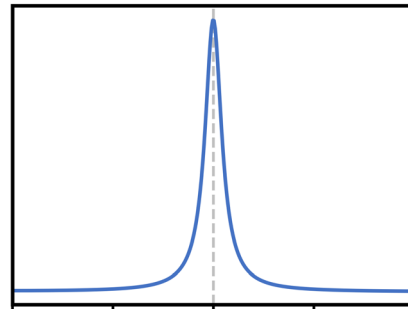
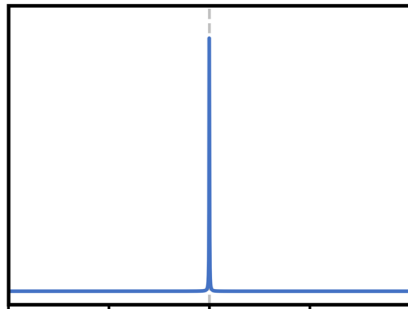
$$\sim \Im \pi$$

$$\sim \Re \pi$$

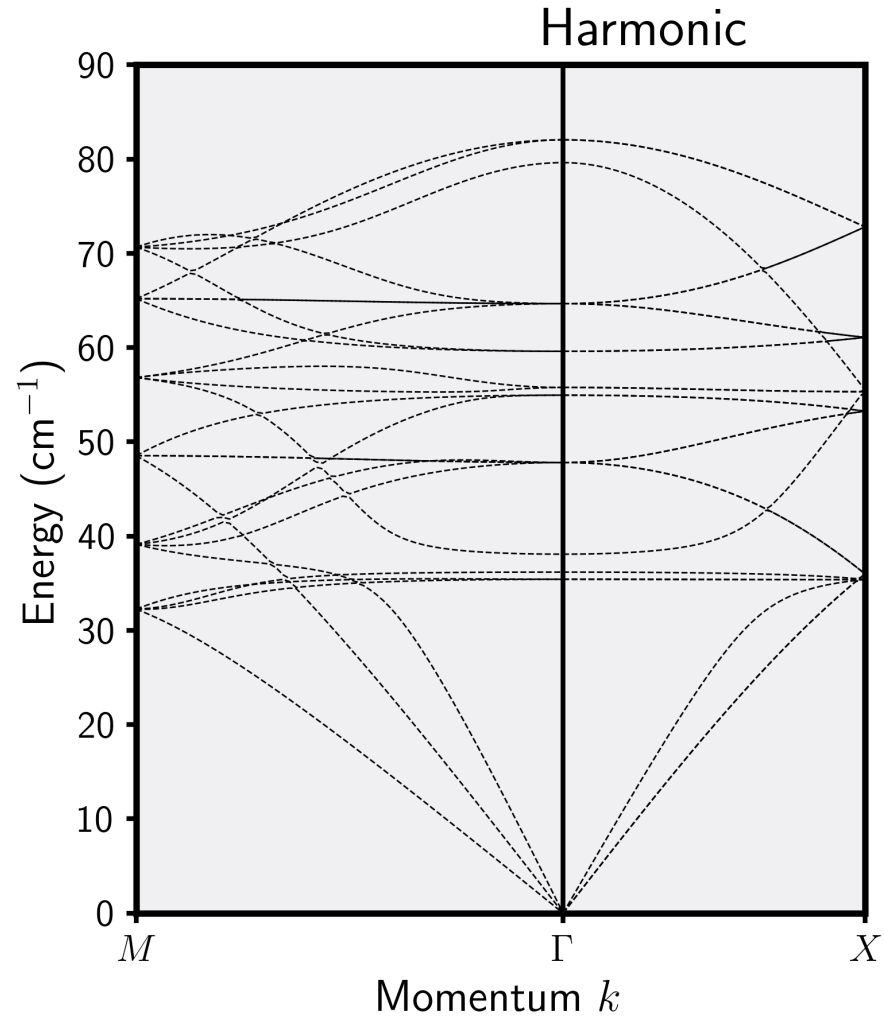
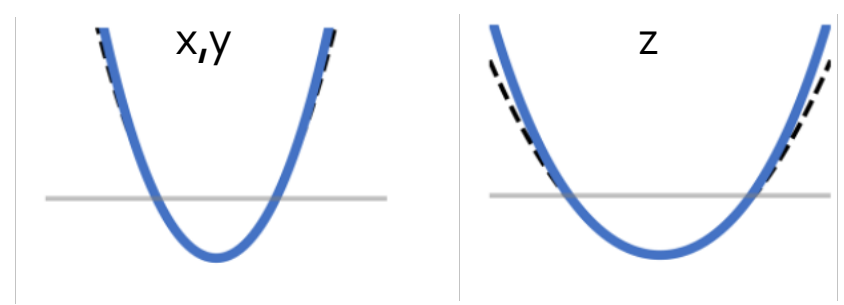
time  
 $\langle u(t)u(0) \rangle$



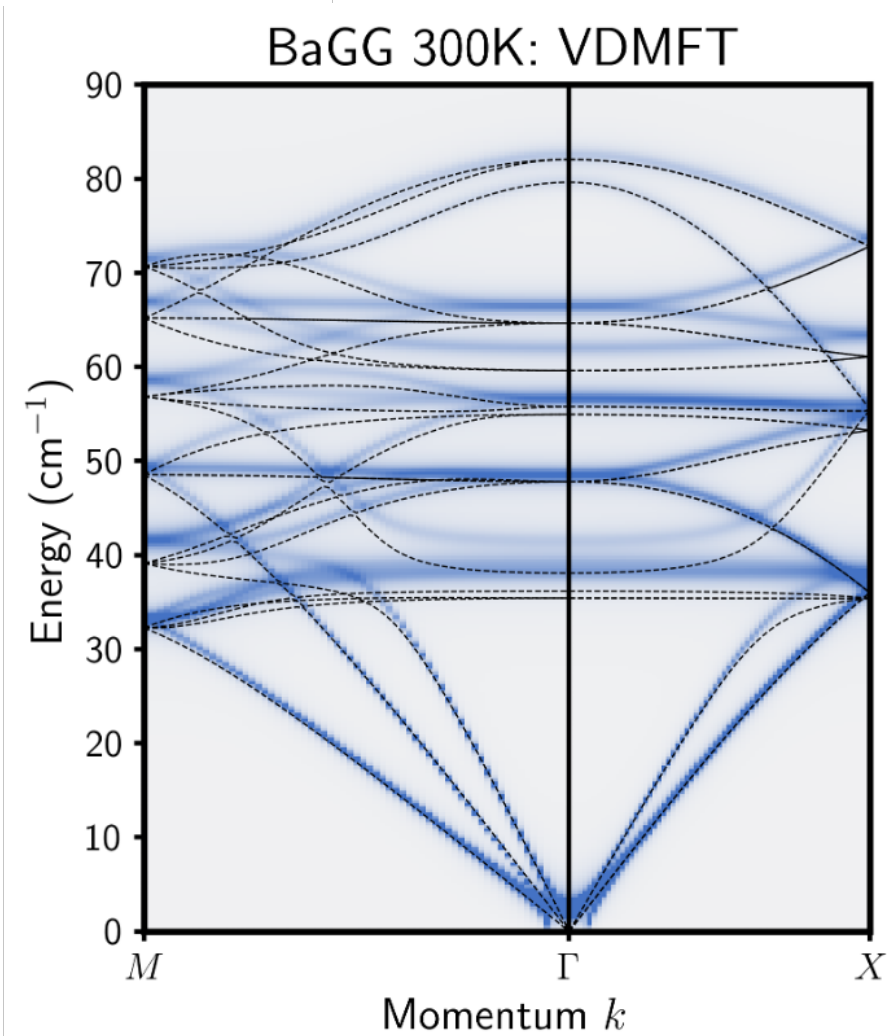
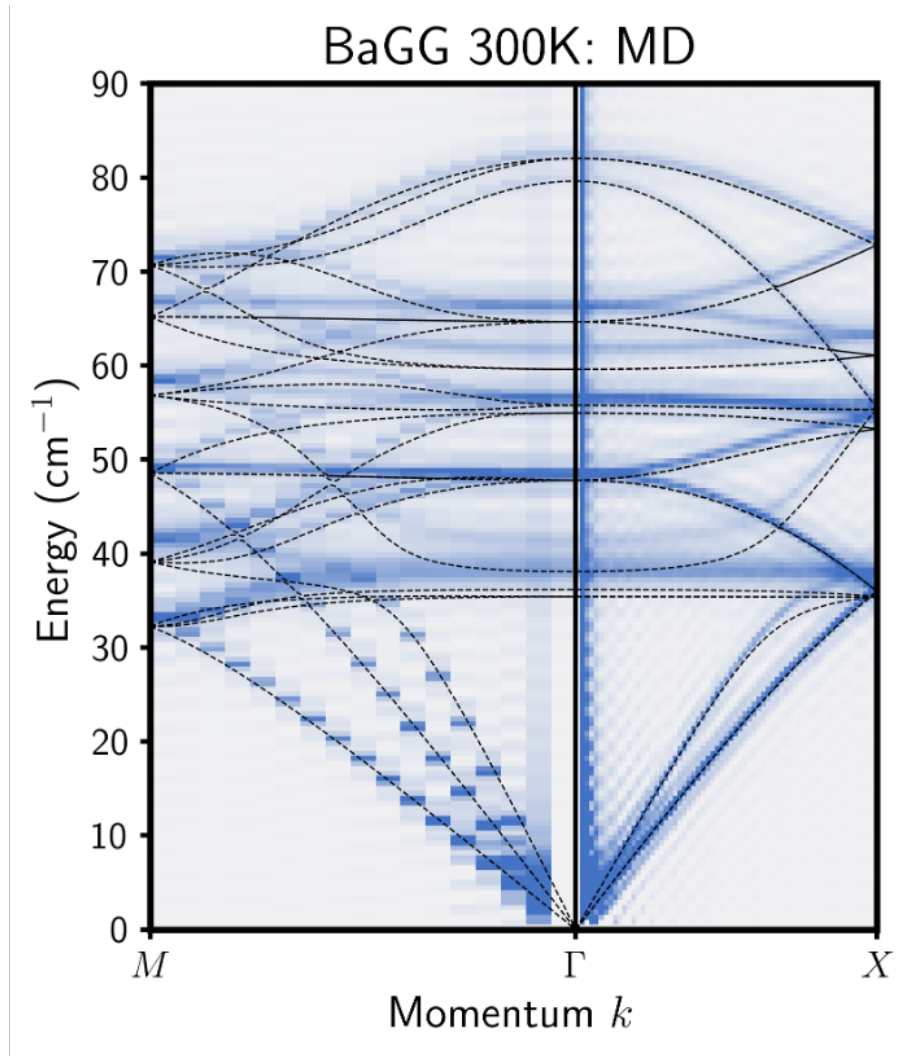
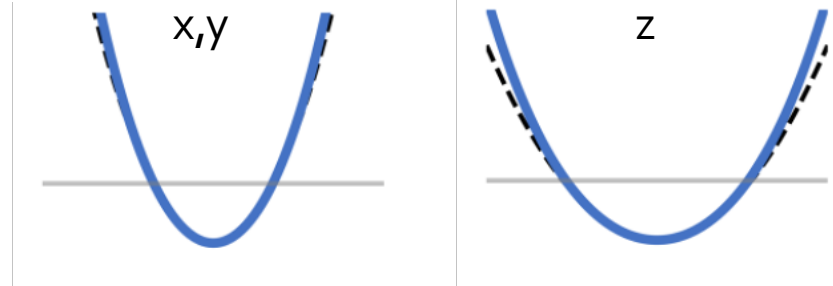
frequency  
 $A(\omega)$



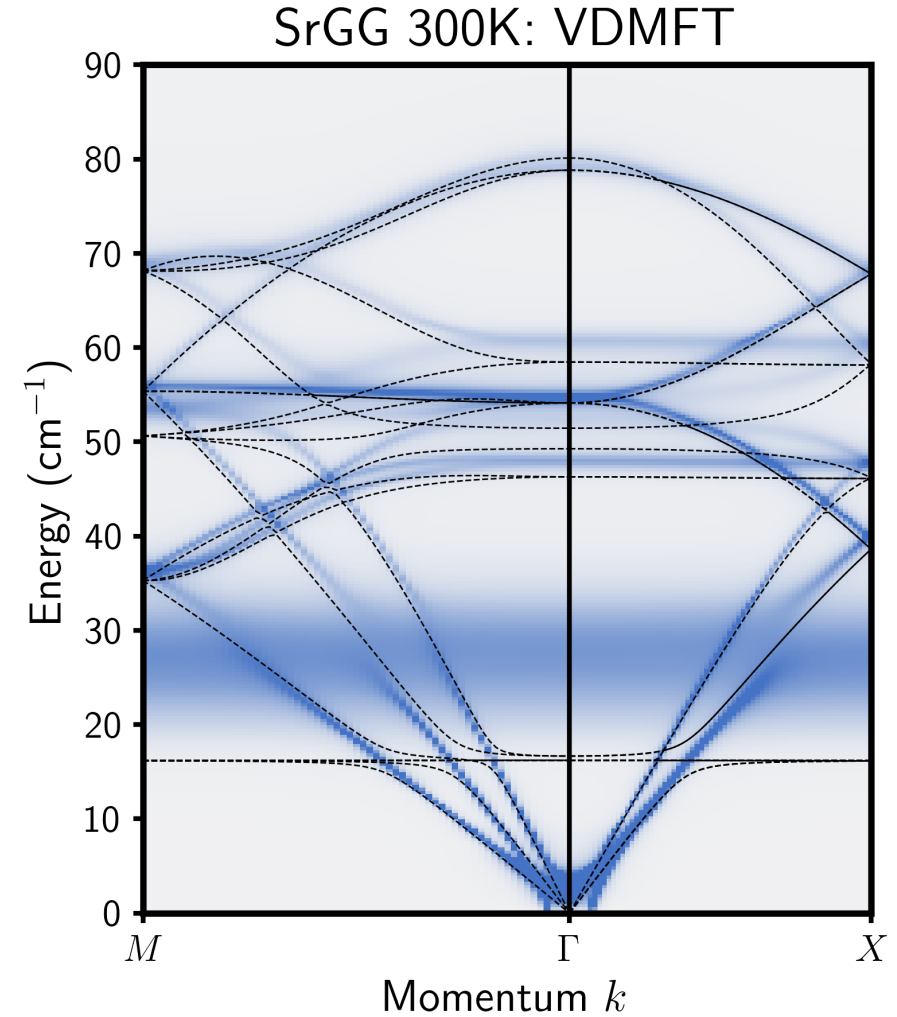
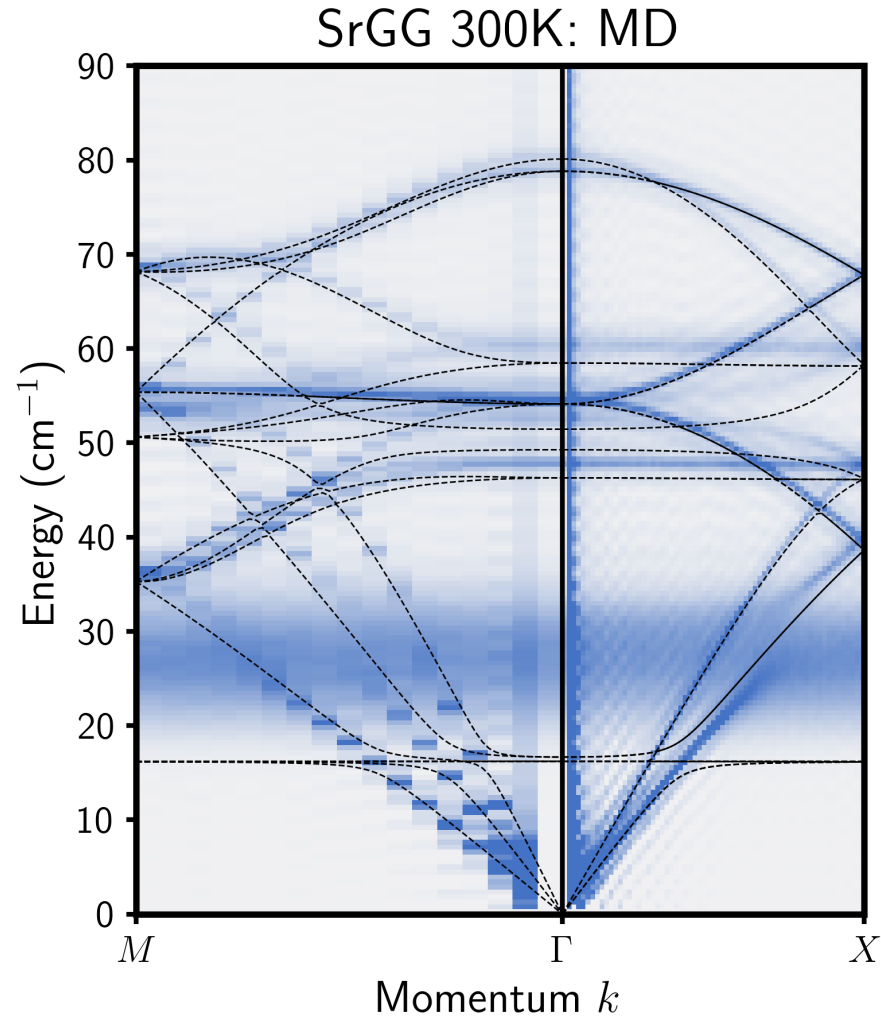
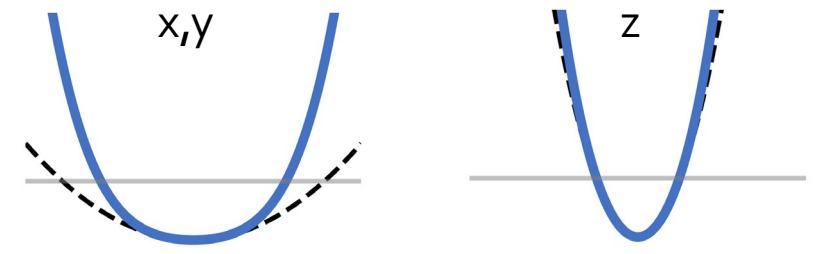
# VDMFT: BaGG at 300K



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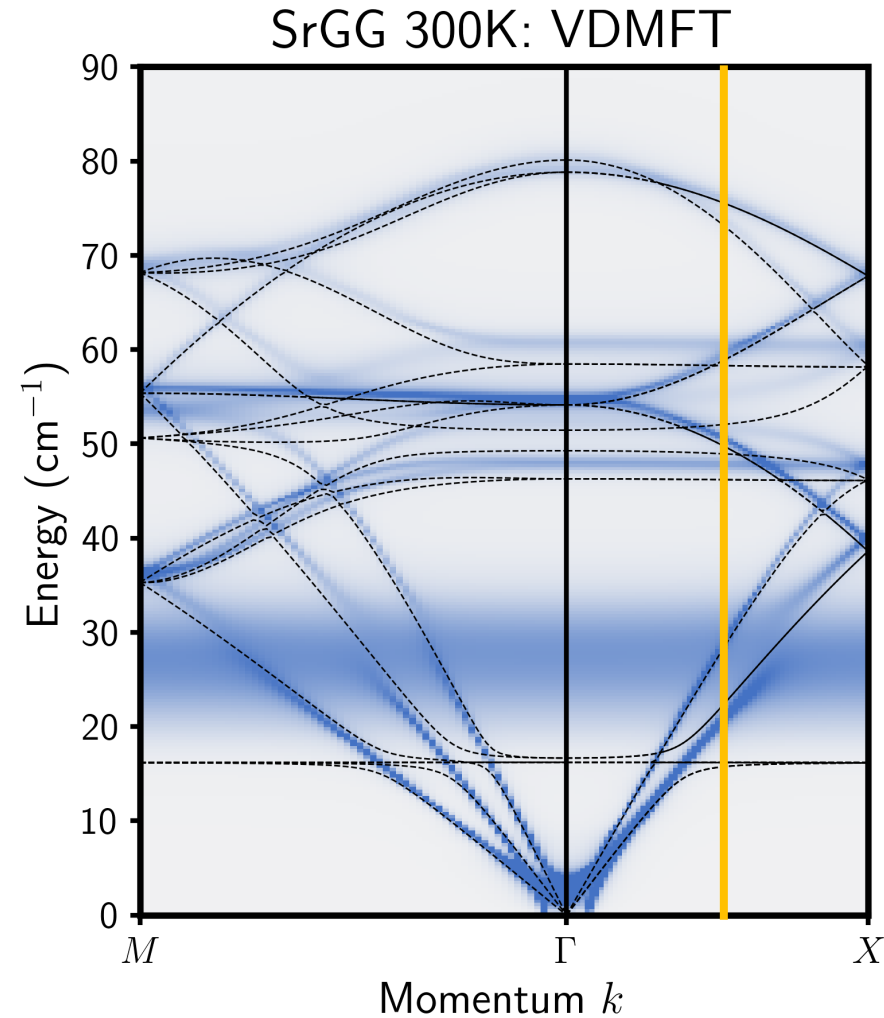
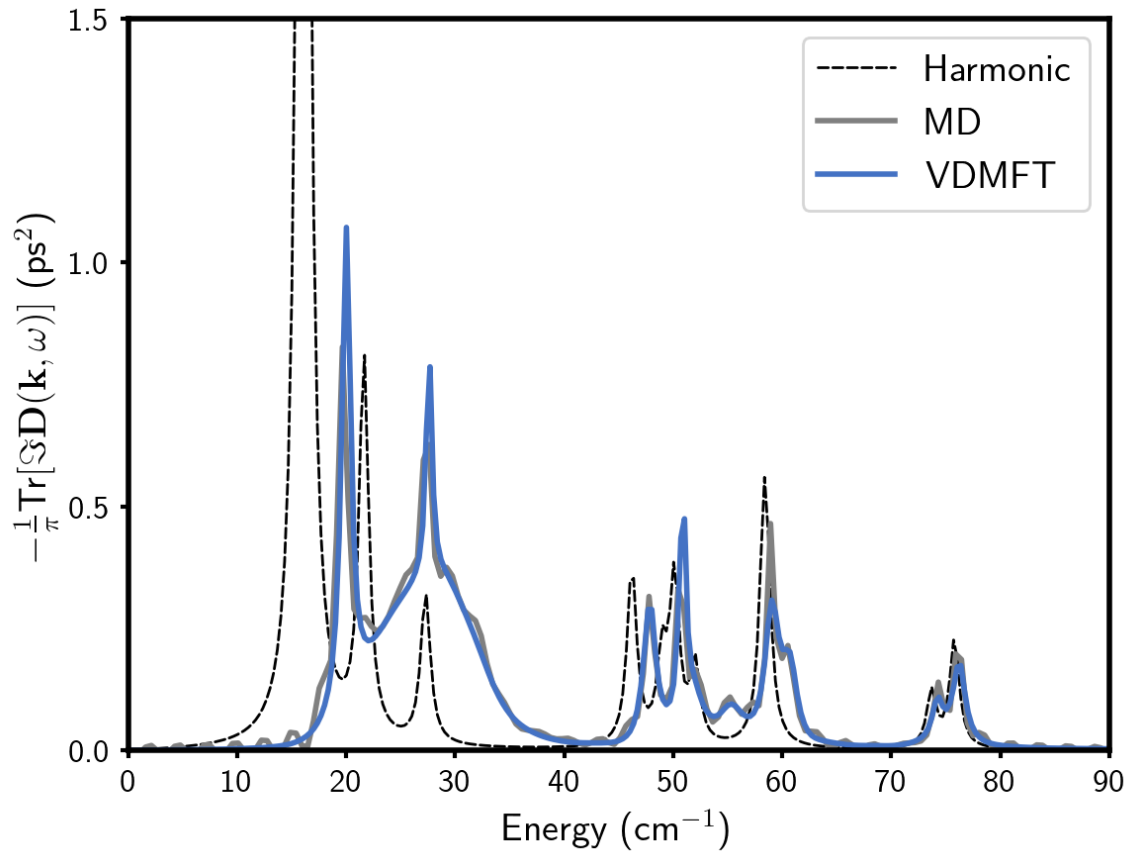
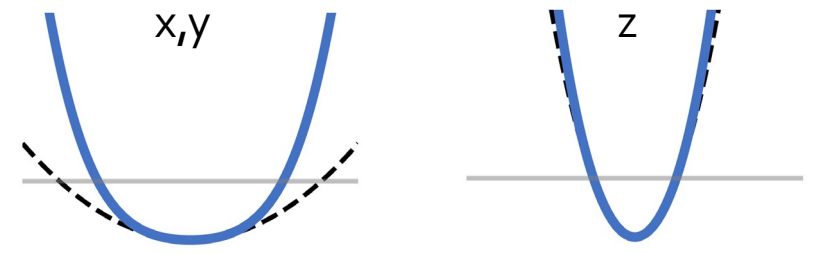


# Turning up anharmonicity: SrGG at 300K

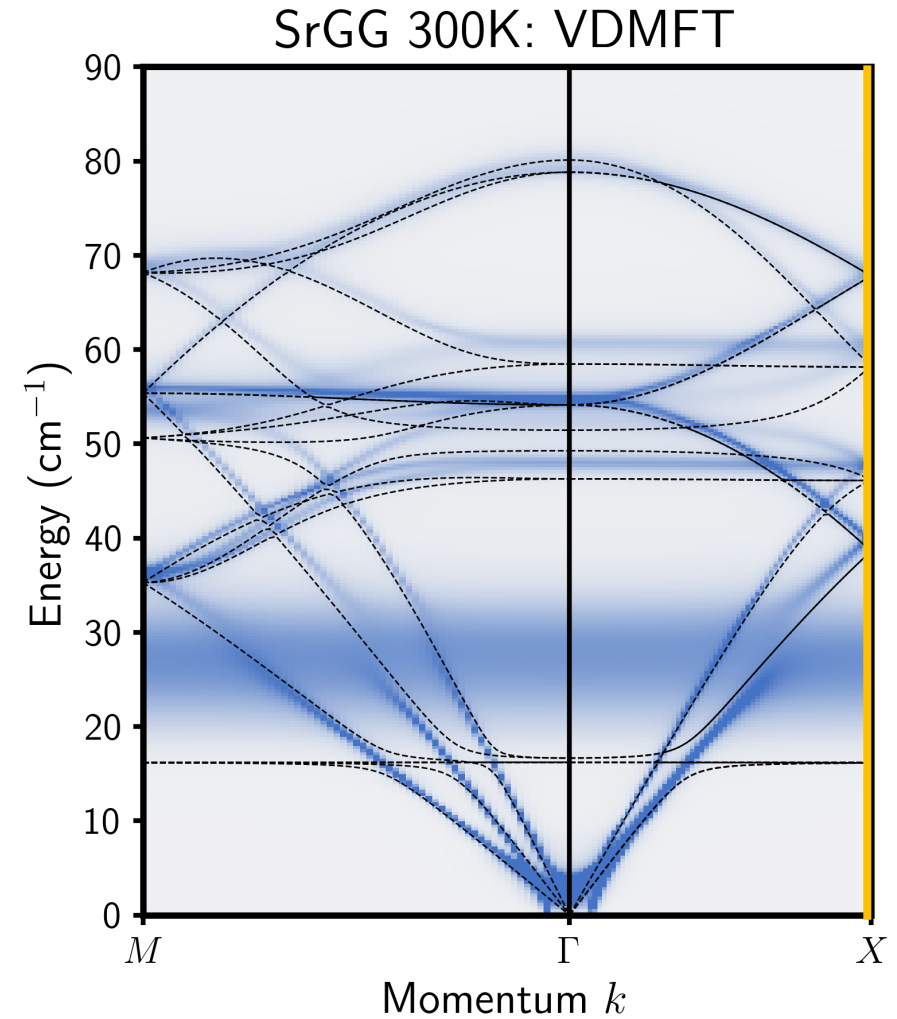
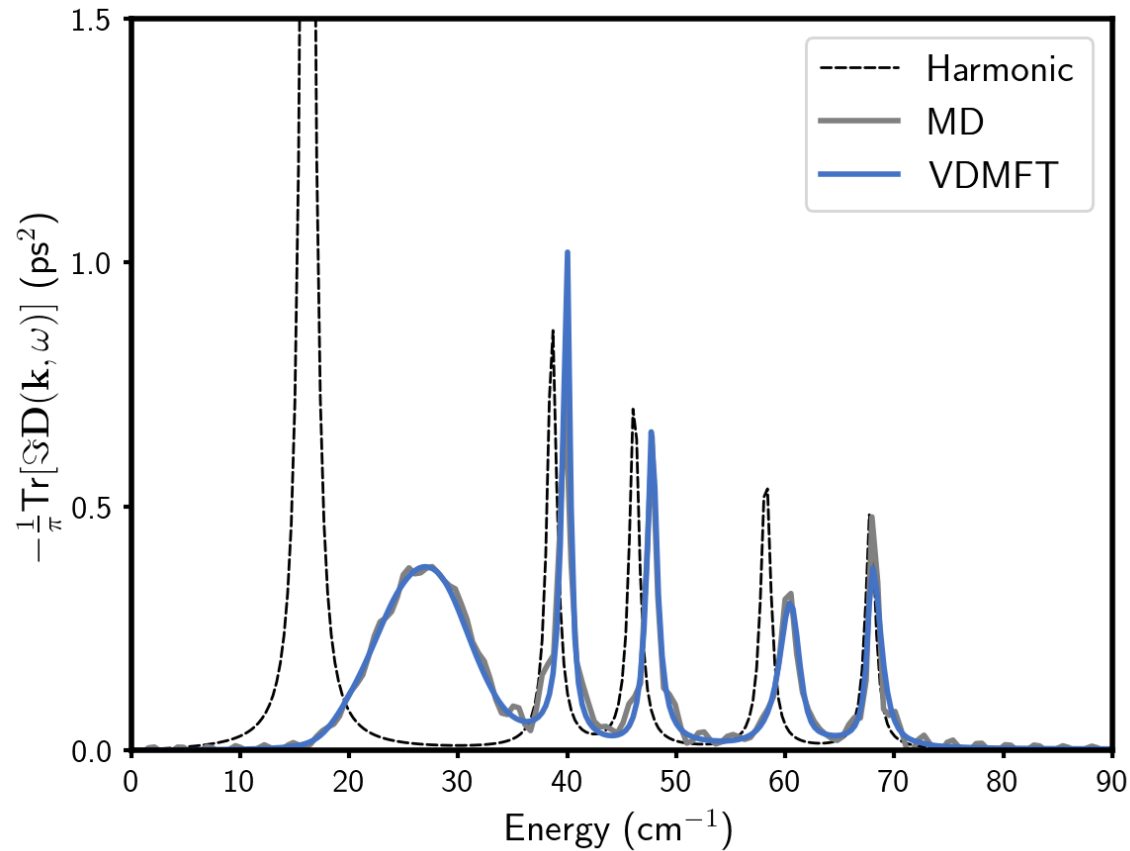
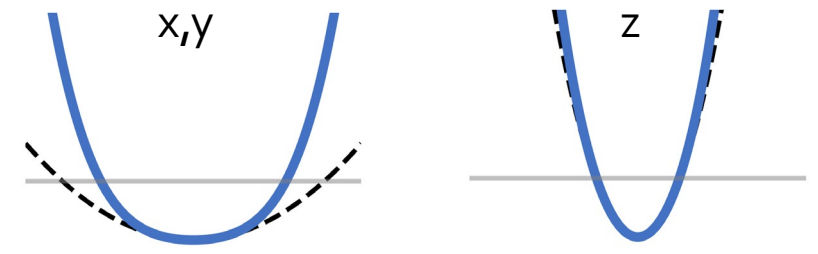




# Turning up anharmonicity: SrGG at 300K

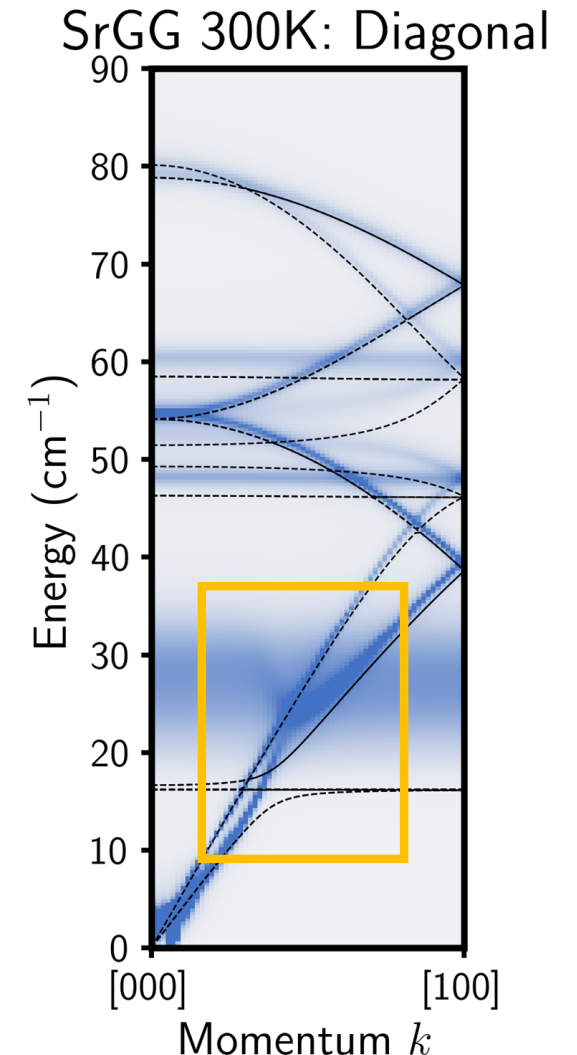
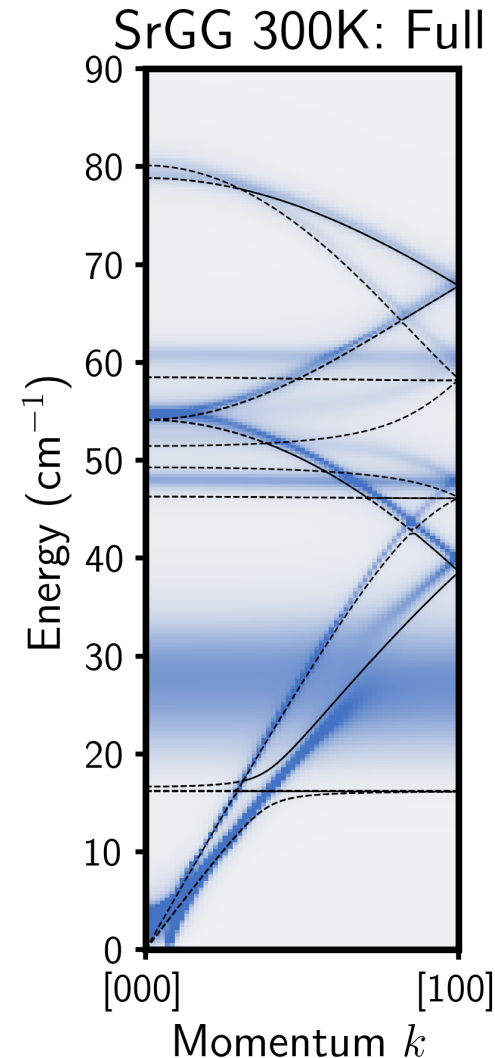


# Turning up anharmonicity: SrGG at 300K



# How important is the *full* self-energy?

- Diagonal approximation
  - Neglects non-diagonal elements of self-energy:  $D^{-1} \rightarrow D$
- Significant **anharmonic** hybridization between cage-acoustic and guest-rattling modes
- Phonon picture is not valid!

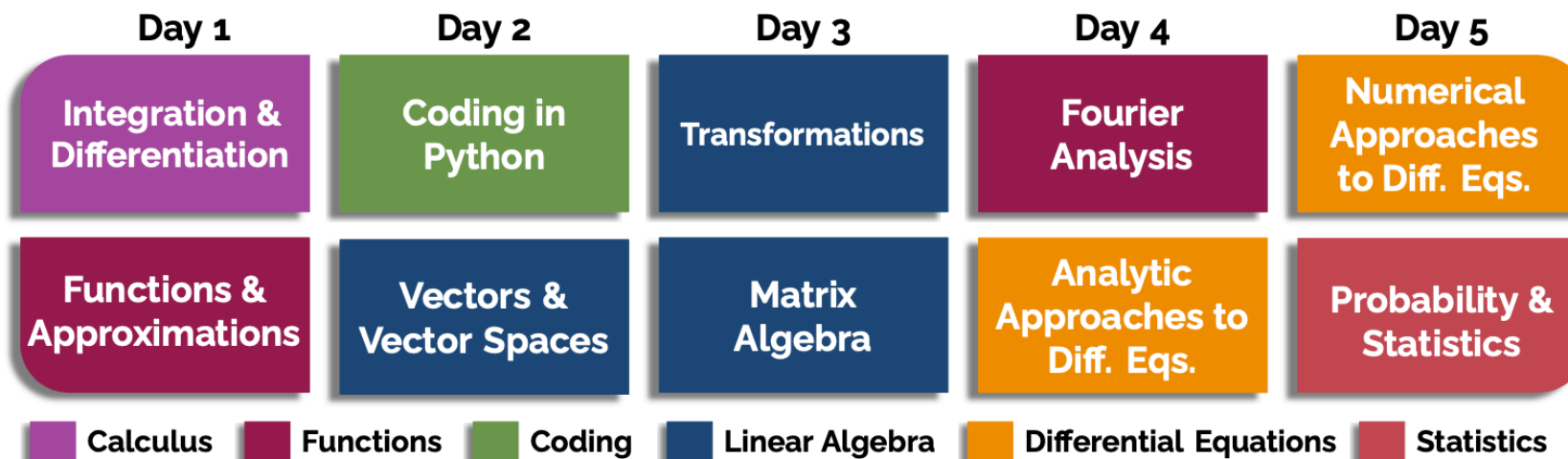


# Summary and looking forward

- VDMFT calculates anharmonic phonon GF
  - Includes local anharmonicity exactly
  - Successfully calculates phonon spectral function for model clathrates with various degrees of anharmonicity
  - Cheaper than MD and can account for nuclear quantum effects
- Can use phonon GF to compute thermal conductivity
  - Phonon hybridization is key to clathrates' low thermal conductivity

# College of Chemistry Math Bootcamp

- Student-initiated, designed, and taught for incoming physical chemistry students centered around group problem solving



- Content freely available at [chemmathbootcamp.com](https://chemmathbootcamp.com)

# Acknowledgements




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