

Computing Bifurcation & Stability Properties of Crystals

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- Introduction & motivation
 - Active materials & Martensitic transformations (MT's)
 - Shape Memory Alloys (SMA's)
- Atomistic modeling of MT's
 - Temperature-dependent atomic potentials
 - Bifurcation & stability investigation of stress-free phases
 - Hysteretic proper MT between cubic $B2$ and orthorhombic $B19$ phases
- Computational challenges
 - Crystal stability
 - Equilibrium path following
 - Behavior near bifurcation points
- Summary & conclusions

Active Materials

- Multi-physics coupling — Crystal structure changes

Magnetostrictive Materials

Magnetic Field



Mechanical

Ferroelectric Materials

Electric Field



Mechanical

Shape Memory Alloys

Temperature

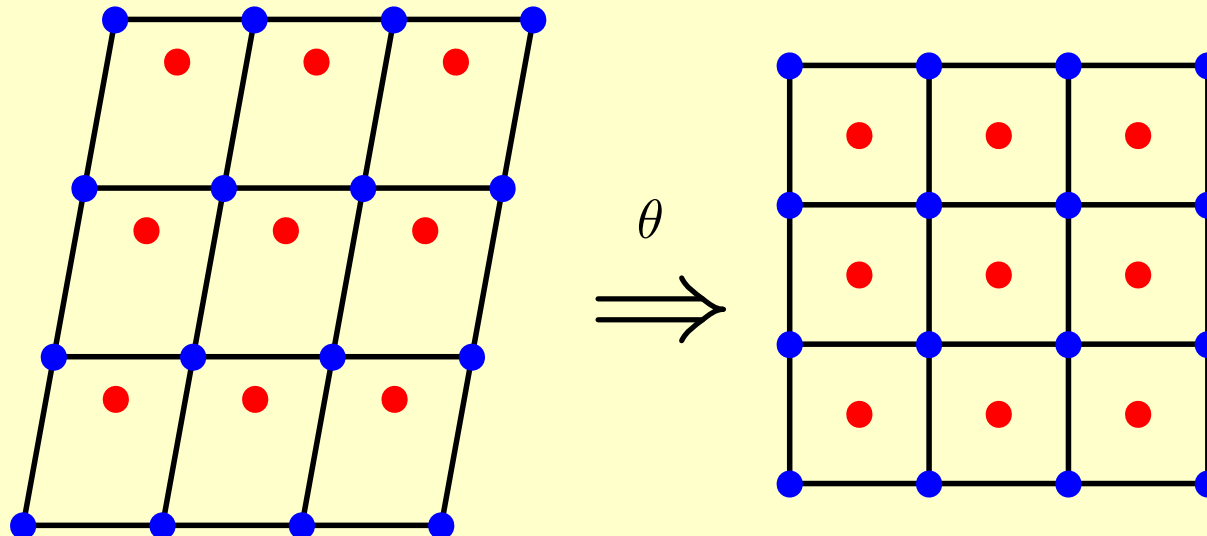


Mechanical

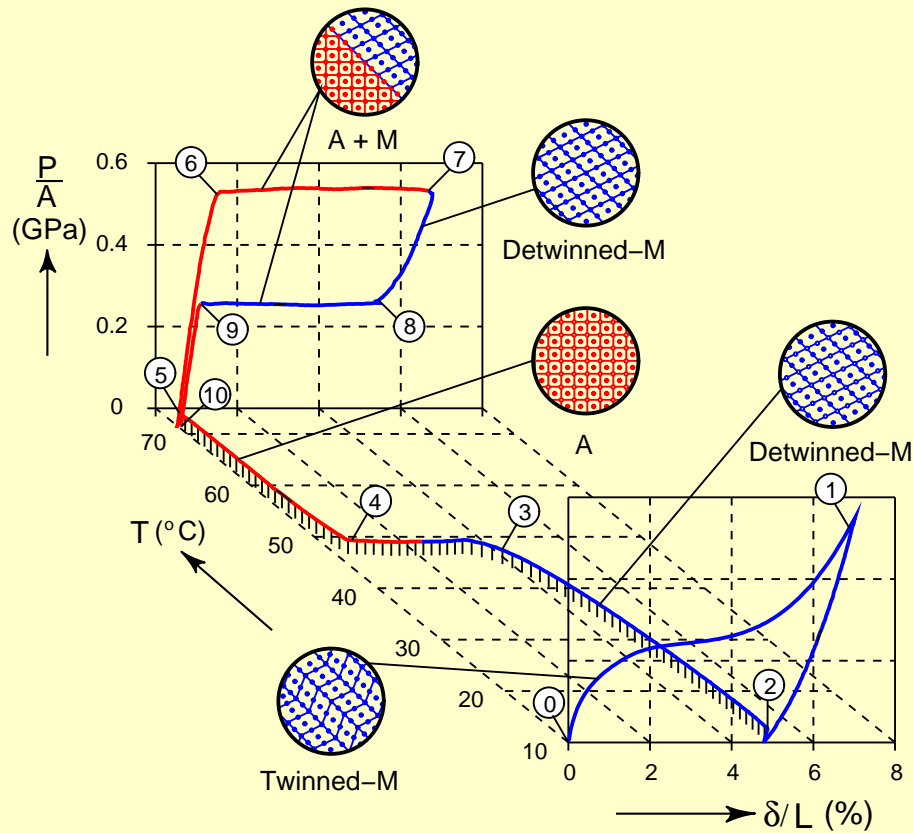
Martensitic Transformations
Materials on the cusp of an instability

Martensite

Austenite



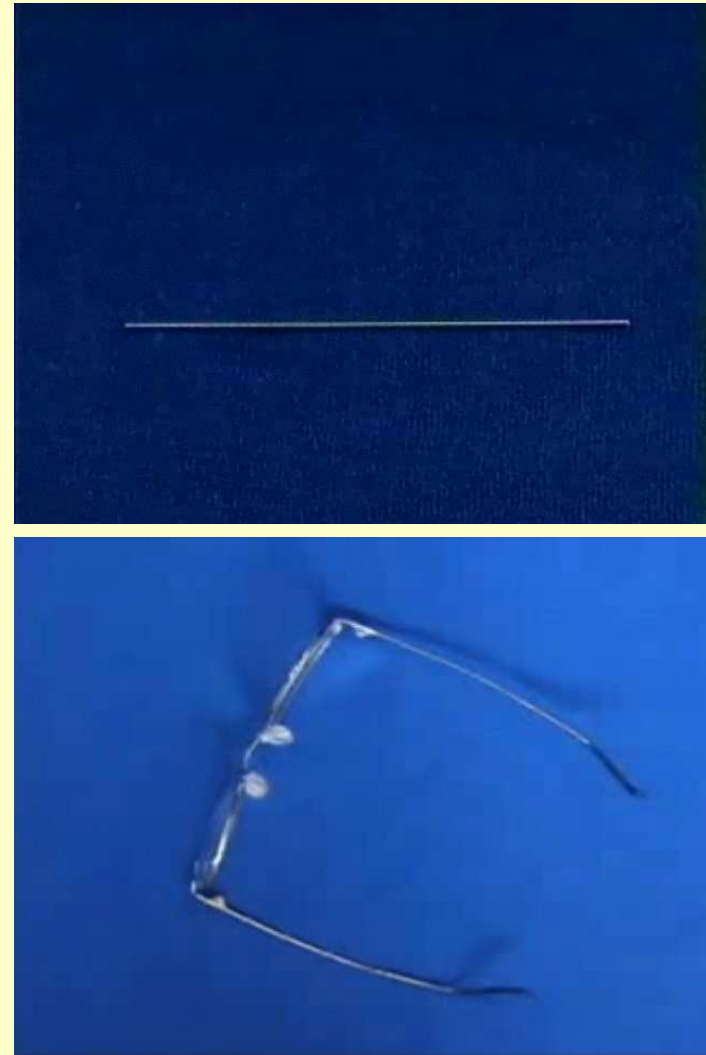
Shape Memory Alloys (SMAs)



Tensile behavior of NiTi

(exhibiting the shape memory effect and pseudo-elasticity)

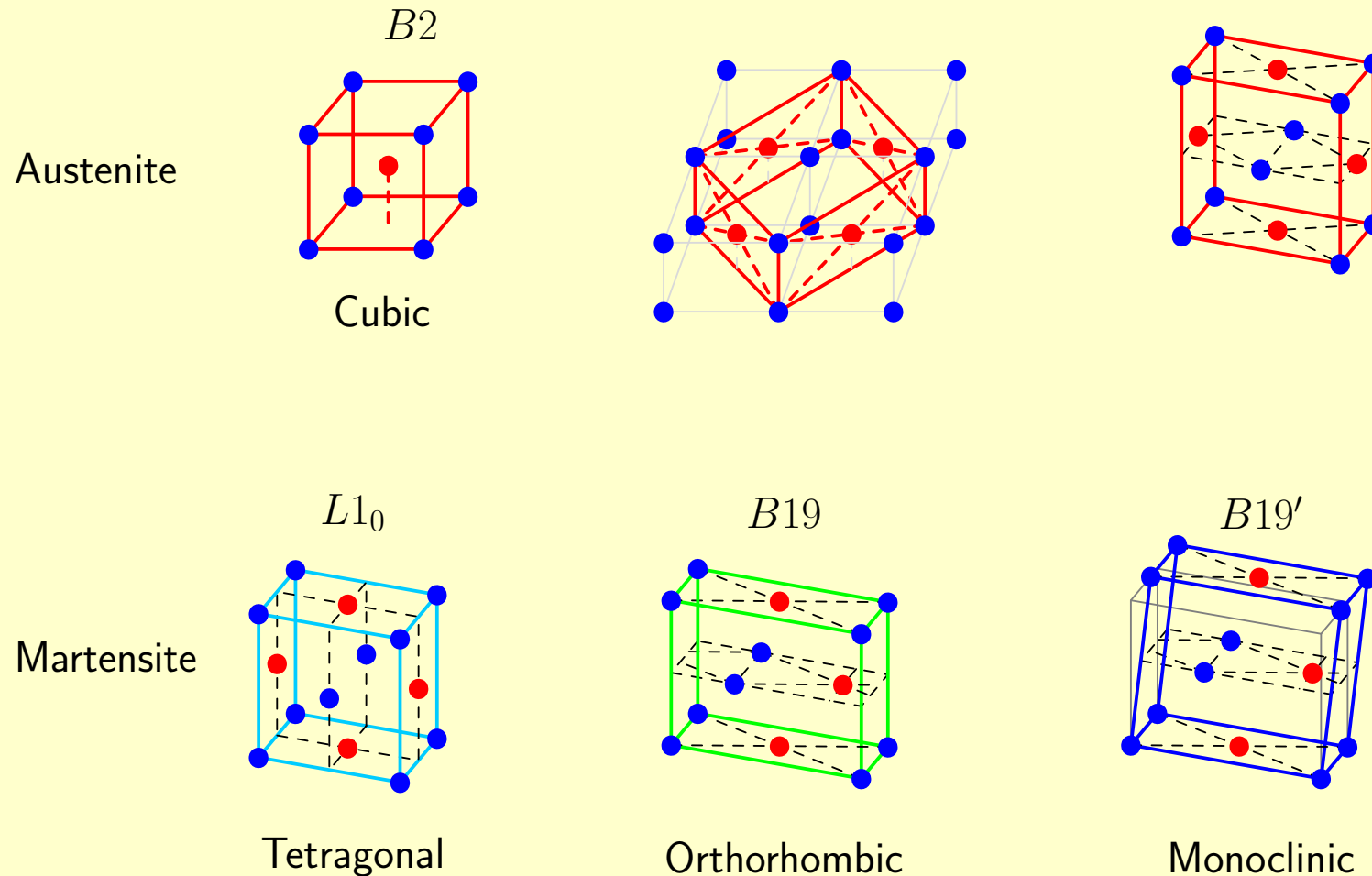
(J. Shaw 1997)



Materials Research Science and Engineering Center
at the University of Wisconsin - Madison
www.mrsec.wisc.edu/nano

The Crystal Structures of SMAs

- Prevalent austenite and martensite crystals in shape memory alloys



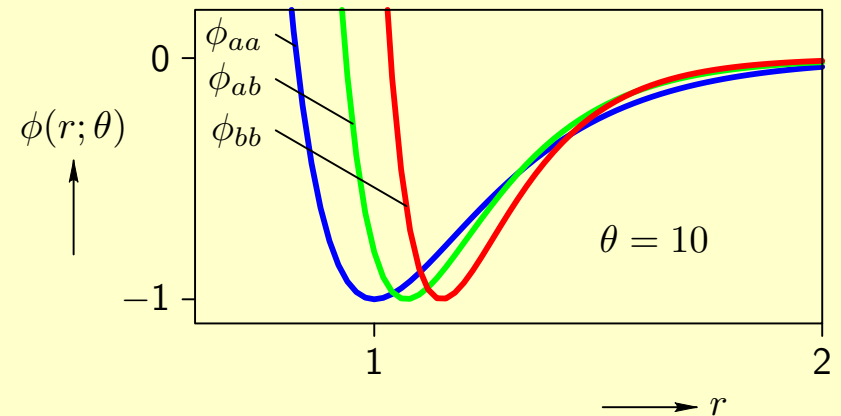
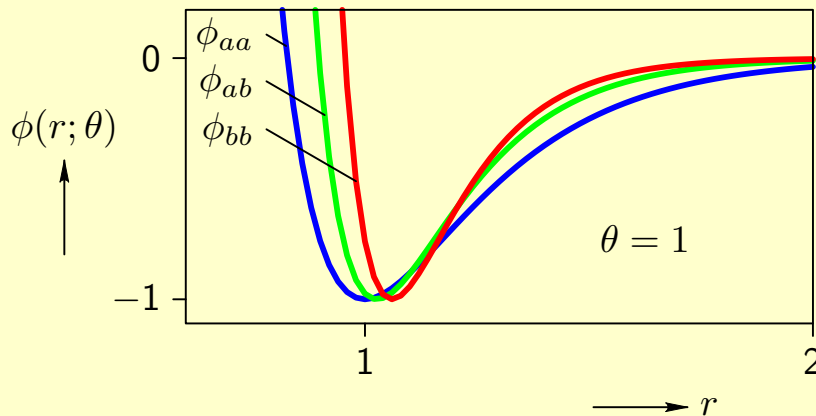
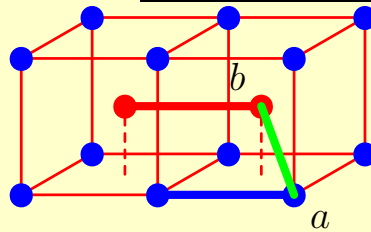
Objective: Develop an atomic model to capture *proper* martensitic transformations such as those found in shape memory alloys

Pair-Potential Model

$$\phi(r; \theta) = A \left\{ \exp \left[-2B \left(\frac{r}{\hat{r}(\theta)} - 1 \right) \right] - 2 \exp \left[-B \left(\frac{r}{\hat{r}(\theta)} - 1 \right) \right] \right\}$$

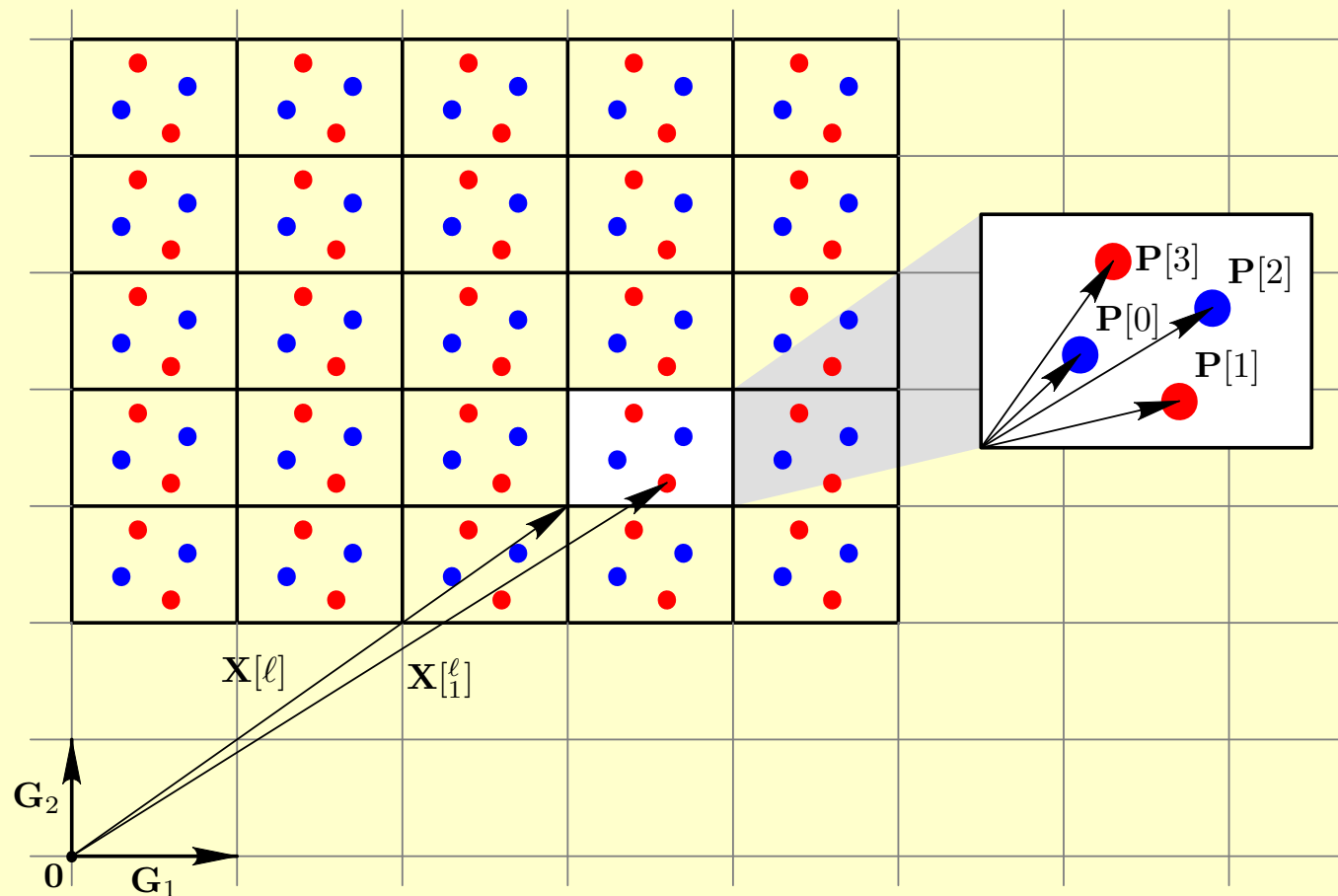
$$\hat{r}(\theta) = r_0 + r_\theta (\theta - 1)$$

	r_0	r_θ	β	A	mass
aa	1	0	4	1	1
bb	1.060	0.010	7	1	0.816
ab	1.026	0.005	5.5	1	N/A



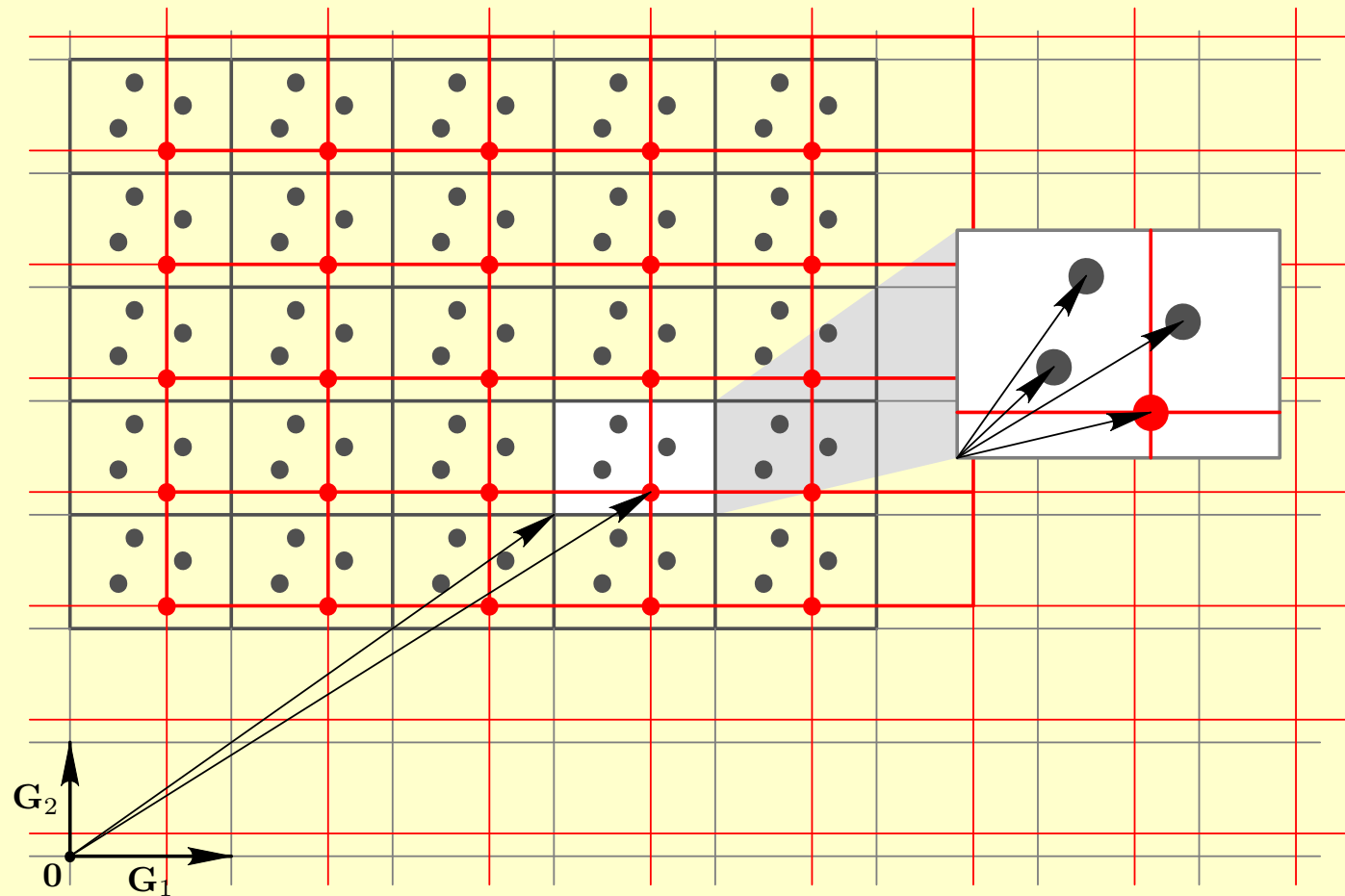
Multilattice (Cauchy-Born Kinematics)

- \mathbf{G}_i – ref. lattice basis
- $\mathbf{X}[\ell]$ – unit-cell ref. pos.
- $\mathbf{X}[\ell_\alpha]$ – reference pos.
- $\mathbf{P}[\alpha]$ – fractional pos.
 $\alpha = 0, 1, 2, 3$



Multilattice (Cauchy-Born Kinematics)

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Multilattice (Cauchy-Born Kinematics)

\mathbf{G}_i – ref. lattice basis

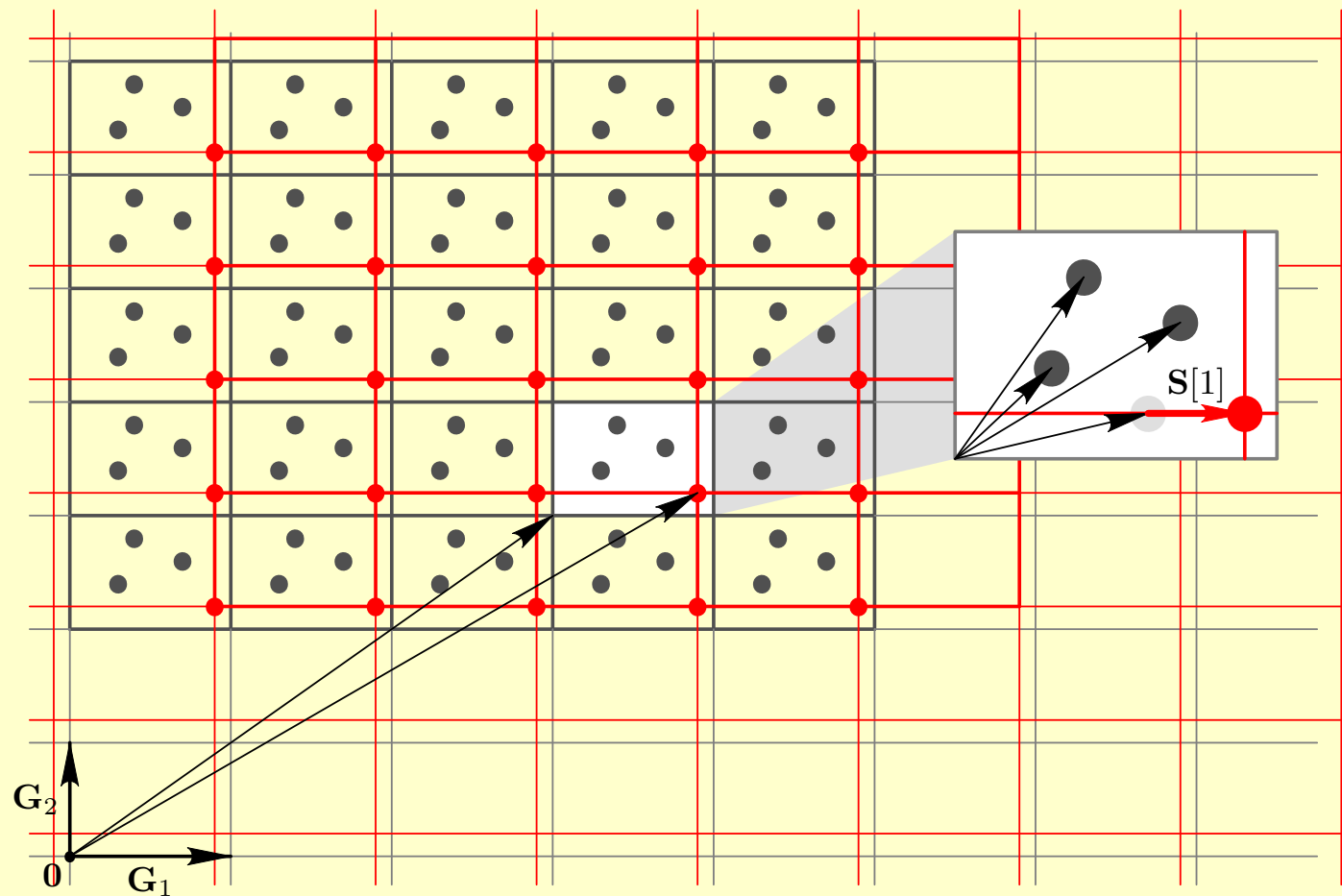
$\mathbf{X}[\ell]$ – unit-cell ref. pos.

$\mathbf{X}[\alpha]$ – reference pos.

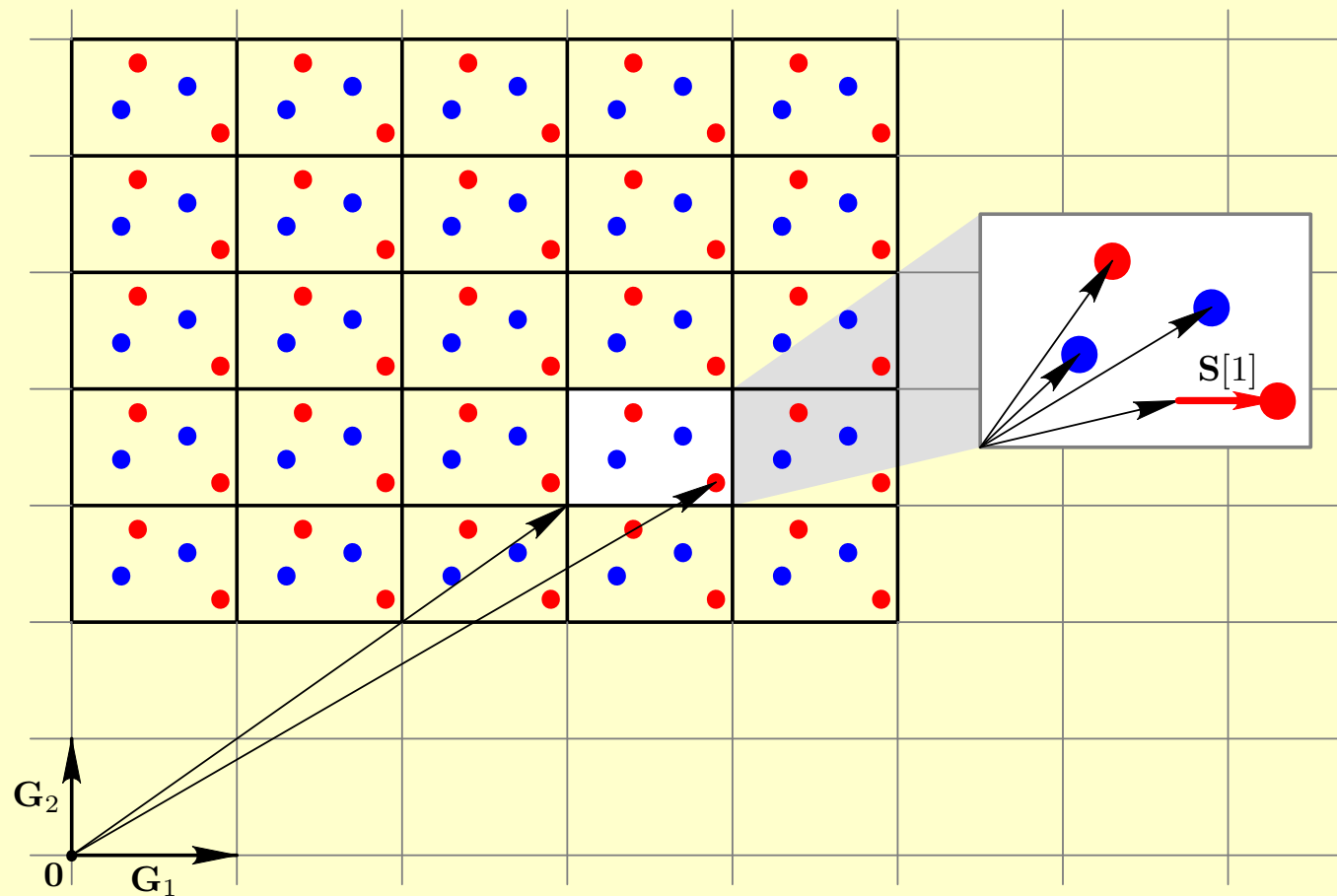
$\mathbf{P}[\alpha]$ – fractional pos.

$$\alpha = 0, 1, 2, 3$$

$\mathbf{S}[\alpha]$ – sub-lat. ref. shifts



Multilattice (Cauchy-Born Kinematics)



G_i – ref. lattice basis

$X[\ell]$ – unit-cell ref. pos.

$X[\alpha^\ell]$ – reference pos.

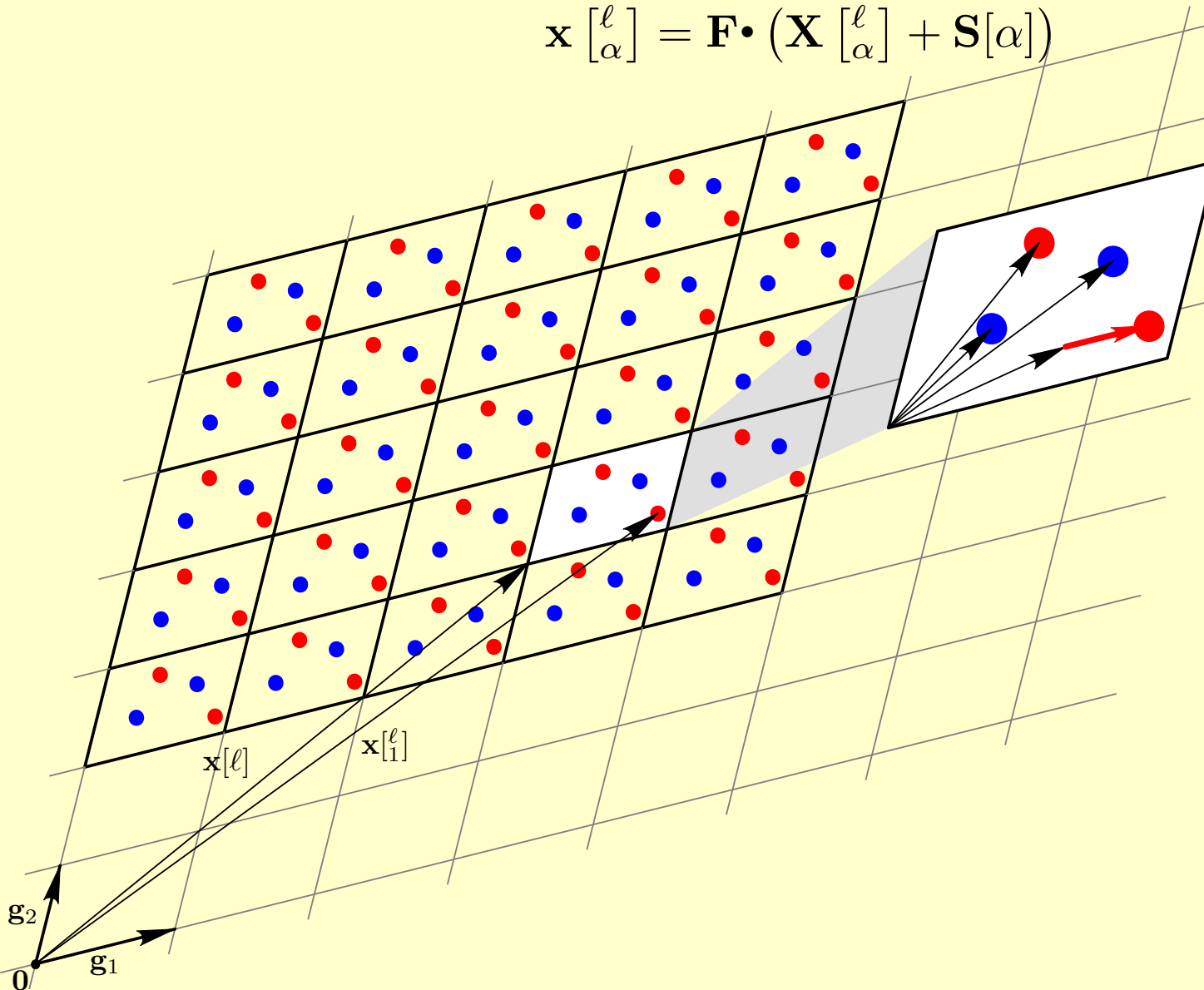
$P[\alpha]$ – fractional pos.

$$\alpha = 0, 1, 2, 3$$

$S[\alpha]$ – sub-lat. ref. shifts

Multilattice (Cauchy-Born Kinematics)

$$\mathbf{x} \begin{bmatrix} \ell \\ \alpha \end{bmatrix} = \mathbf{F} \cdot (\mathbf{X} \begin{bmatrix} \ell \\ \alpha \end{bmatrix} + \mathbf{S}[\alpha])$$



\mathbf{G}_i – ref. lattice basis

$\mathbf{X}[\ell]$ – unit-cell ref. pos.

$\mathbf{X} \begin{bmatrix} \ell \\ \alpha \end{bmatrix}$ – reference pos.

$\mathbf{P}[\alpha]$ – fractional pos.

$$\alpha = 0, 1, 2, 3$$

$\mathbf{S}[\alpha]$ – sub-lat. ref. shifts

\mathbf{F} – uniform deformation

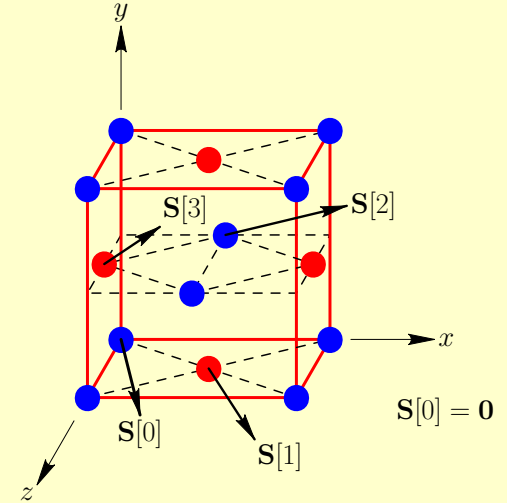
\mathbf{g}_i – current lattice basis

$\mathbf{x}[\ell]$ – unit-cell current pos.

$\mathbf{x} \begin{bmatrix} \ell \\ \alpha \end{bmatrix}$ – current pos.

$\mathbf{X} \begin{bmatrix} \ell \\ \alpha \end{bmatrix}$ — reference position vector of atom α in unit cell ℓ

$\mathbf{S}[\alpha]$ — displacement vector of atom α (sub-lattice)
 $\alpha = 0, 1, 2, 3$



- Current position vector (Cauchy-Born kinematics, $\alpha = 0, 1, 2, 3$)

$$\mathbf{x} \begin{bmatrix} \ell \\ \alpha \end{bmatrix} = \mathbf{F} \cdot (\mathbf{X} \begin{bmatrix} \ell \\ \alpha \end{bmatrix} + \mathbf{S}[\alpha])$$

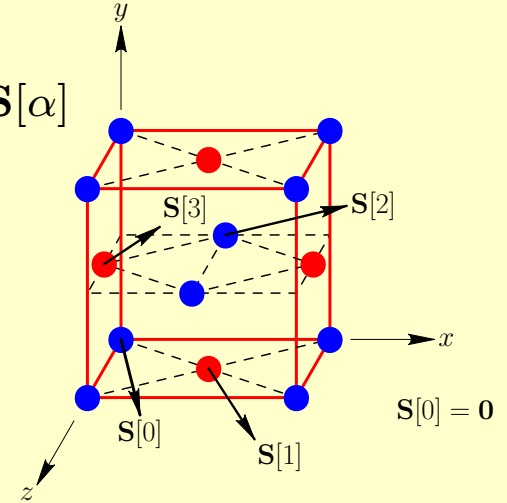
- Energy density

$$\tilde{W}(\mathbf{u}; \theta) = \frac{1}{2V} \sum_{\alpha'} \sum_{\begin{bmatrix} \ell \\ \alpha \end{bmatrix}} \phi_{\alpha\alpha'} (r \begin{bmatrix} \ell & 0 \\ \alpha & \alpha' \end{bmatrix}; \theta)$$

$$\mathbf{u} \equiv \{\mathbf{F}, \mathbf{S}[1], \mathbf{S}[2], \mathbf{S}[3]\}, \quad r \begin{bmatrix} \ell & \ell' \\ \alpha & \alpha' \end{bmatrix} = \left\| \mathbf{x} \begin{bmatrix} \ell \\ \alpha \end{bmatrix} - \mathbf{x} \begin{bmatrix} \ell' \\ \alpha' \end{bmatrix} \right\|$$

- Equilibrium: 15 DOFs — 6 from $\mathbf{U} = \mathbf{U}^T$ and 9 from $\mathbf{S}[\alpha]$

$$\frac{\partial \tilde{W}}{\partial \mathbf{u}} = \mathbf{0} \left\{ \begin{array}{l} \frac{\partial \tilde{W}}{\partial \mathbf{U}} = \mathbf{0}, \\ \frac{\partial \tilde{W}}{\partial \mathbf{S}[1]} = \mathbf{0}, \quad \frac{\partial \tilde{W}}{\partial \mathbf{S}[2]} = \mathbf{0}, \quad \frac{\partial \tilde{W}}{\partial \mathbf{S}[3]} = \mathbf{0}. \end{array} \right.$$



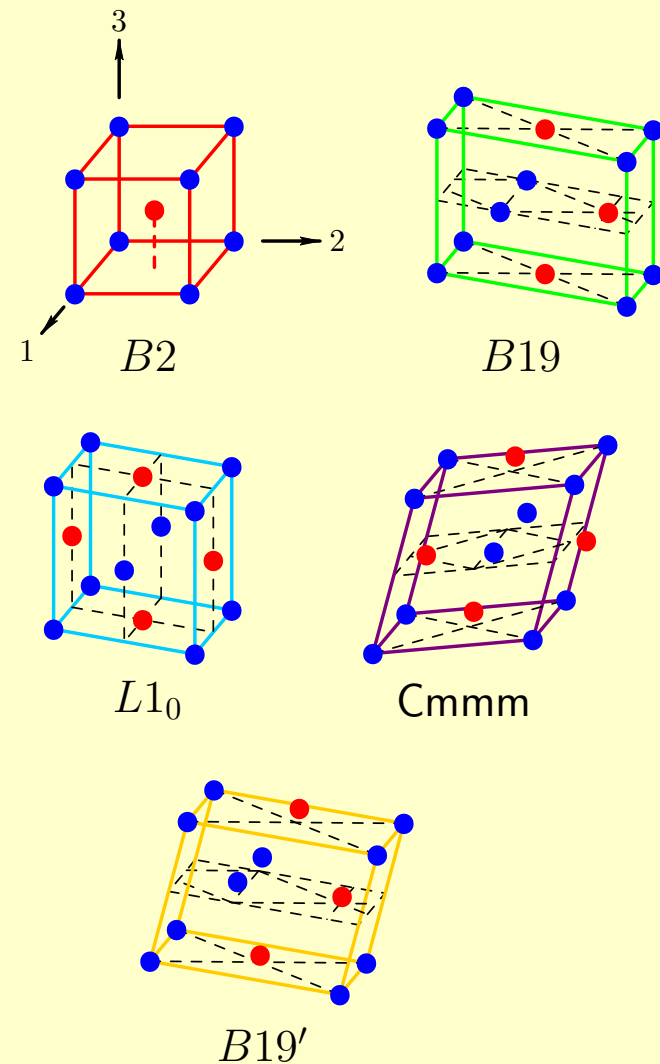
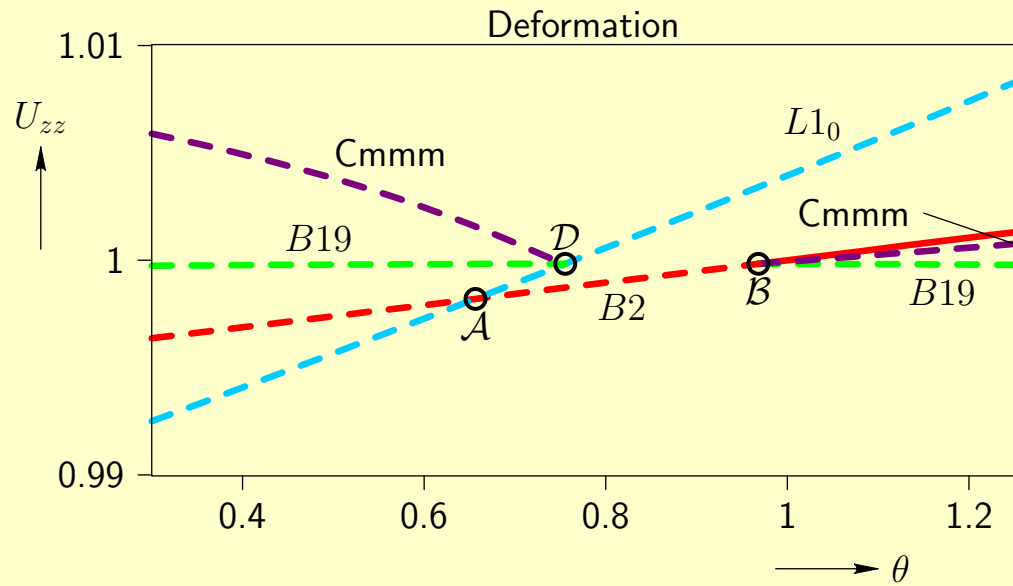
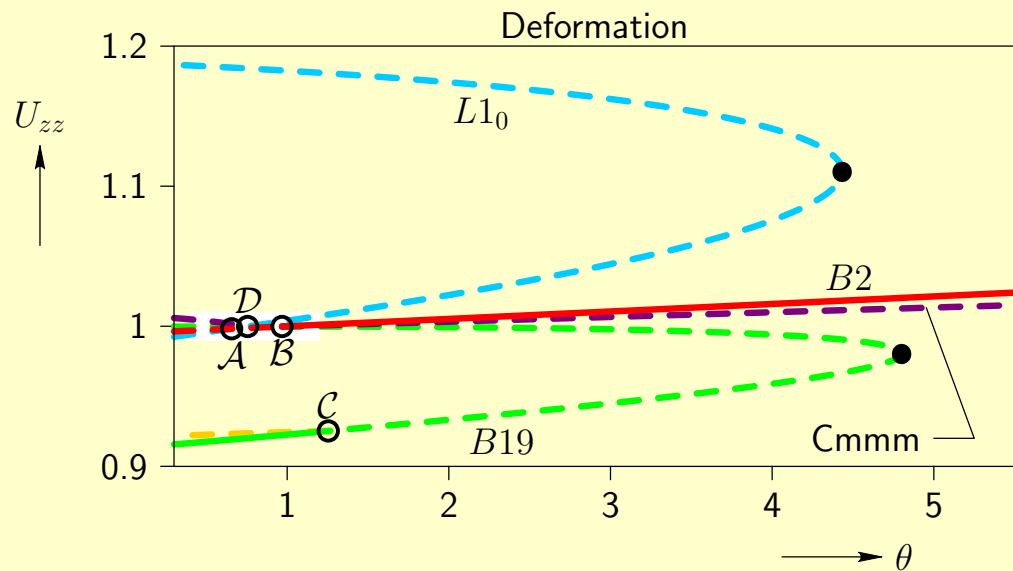
- Stability
 - Cauchy-Born stability (local energy minimizer):

$$\delta \mathbf{u} \frac{\partial^2 \tilde{W}}{\partial \mathbf{u} \partial \mathbf{u}} \delta \mathbf{u} > 0; \quad \delta \mathbf{u} = \{\delta \mathbf{U}, \delta \mathbf{S}[1], \delta \mathbf{S}[2], \delta \mathbf{S}[3]\}, \quad \delta \mathbf{U} = \delta \mathbf{U}^T.$$

- Phonon stability:

$$\left(\omega^{(q)}(\mathbf{k}) \right)^2 > 0, \quad \forall \mathbf{k}, q.$$

4-Lattice Bifurcation Diagram



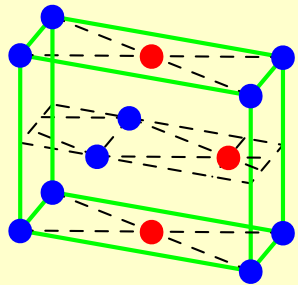
● Hysteretic proper Martensitic transformation between $B2$ & $B19$

Transformation Parameters

$$B2 \implies B19$$

Martensitic Transformation

- Experimental right stretch tensor



$B19$

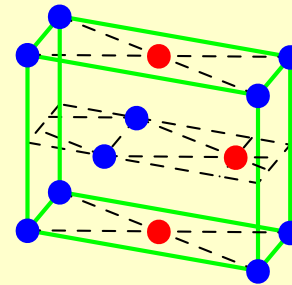
$$\mathbf{U} = \begin{bmatrix} 1.024 & 0.0106 & 0 \\ 0.0106 & 1.024 & 0 \\ 0 & 0 & 0.9491 \end{bmatrix}$$

AuCd, (*Chang, Read (1951)*)

$$\mathbf{U} = \begin{bmatrix} 1.042 & 0.0194 & 0 \\ 0.0194 & 1.042 & 0 \\ 0 & 0 & 0.9178 \end{bmatrix}$$

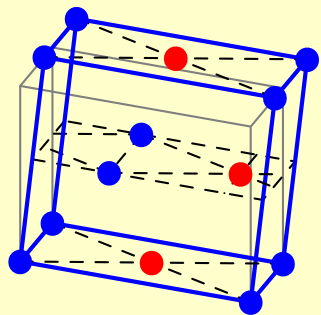
CuAlNi, (*Otsuka, Shimizu (1974)*)

- Simulated right stretch tensor ($\theta = 1.0$)



$B19$

$$\mathbf{U} = \begin{bmatrix} 1.045 & 0.0173 & 0 \\ 0.0173 & 1.045 & 0 \\ 0 & 0 & 0.9224 \end{bmatrix}$$



$B19'$

$$\mathbf{U} = \begin{bmatrix} 1.025 & 0.0620 & 0.0490 \\ 0.0620 & 1.025 & 0.0490 \\ 0.0490 & 0.0490 & 0.9587 \end{bmatrix}$$

NiTi, (*Otsuka et al. (1971)*)

- Crystal structure stability
 - Robust stability criterion — phonon spectra
 - Efficient numerical evaluation of phonon spectra
- Efficient equilibrium path following
 - Reduced set of equations based on symmetry
 - Pseudo-arc-length method
- Determine behavior near bifurcation points
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- Compute all eigenvalues of $K = \frac{\partial^2 \tilde{W}}{\partial \mathbf{u} [\ell_\alpha] \partial \mathbf{u} [\ell_\alpha]}, \quad (3MN \times 3MN)$

M —number of atoms per unit cell

N —number of unit cells in crystal

$$\begin{bmatrix} * & * & * & \cdots \\ * & * & * & \cdots \\ * & * & * & \cdots \\ \vdots & \vdots & \vdots & \ddots \end{bmatrix} \Rightarrow \begin{bmatrix} \lambda_1 & & & \\ & \lambda_2 & & \\ & & \lambda_3 & \\ & & & \ddots \end{bmatrix}$$

- Methods: SVD, LDU, Cholesky, Jacobi, Householder, etc.
- Time complexity: $O([3MN]^3)$. Too slow for large N

Reduction to Block Diagonal Form

- In $3M \times 3M$ block form

$$K = \begin{bmatrix}
 \ddots & & & & & \\
 & \left(K \begin{bmatrix} n-1 & n-1 \\ \alpha & \beta \end{bmatrix} \right) & \left(K \begin{bmatrix} n-1 & n \\ \alpha & \beta \end{bmatrix} \right) & \left(K \begin{bmatrix} n-1 & n+1 \\ \alpha & \beta \end{bmatrix} \right) & & \\
 & \left(K \begin{bmatrix} n & n-1 \\ \alpha & \beta \end{bmatrix} \right) & \left(K \begin{bmatrix} n & n \\ \alpha & \beta \end{bmatrix} \right) & \left(K \begin{bmatrix} n & n+1 \\ \alpha & \beta \end{bmatrix} \right) & & \\
 & \left(K \begin{bmatrix} n+1 & n-1 \\ \alpha & \beta \end{bmatrix} \right) & \left(K \begin{bmatrix} n+1 & n \\ \alpha & \beta \end{bmatrix} \right) & \left(K \begin{bmatrix} n+1 & n+1 \\ \alpha & \beta \end{bmatrix} \right) & & \\
 & & & & \ddots & \\
 & & & & & \ddots
 \end{bmatrix}$$

Reduction to Block Diagonal Form

- In $3M \times 3M$ block form

$$K = \begin{bmatrix} \ddots & & & & & \\ & \left(K \begin{bmatrix} 0 & 0 \\ \alpha & \beta \end{bmatrix} \right) & & & & \\ & \left(K^T \begin{bmatrix} 0 & 1 \\ \alpha & \beta \end{bmatrix} \right) & & & & \\ & \left(K^T \begin{bmatrix} -1 & 1 \\ \alpha & \beta \end{bmatrix} \right) & & & & \\ & & \left(K \begin{bmatrix} 0 & 1 \\ \alpha & \beta \end{bmatrix} \right) & & & \\ & & \left(K \begin{bmatrix} 0 & 0 \\ \alpha & \beta \end{bmatrix} \right) & & & \\ & & \left(K^T \begin{bmatrix} 0 & 1 \\ \alpha & \beta \end{bmatrix} \right) & & & \\ & & & \left(K \begin{bmatrix} -1 & 1 \\ \alpha & \beta \end{bmatrix} \right) & & \\ & & & \left(K \begin{bmatrix} 0 & 1 \\ \alpha & \beta \end{bmatrix} \right) & & \\ & & & \left(K \begin{bmatrix} 0 & 0 \\ \alpha & \beta \end{bmatrix} \right) & & \\ & & & & \ddots & \end{bmatrix}$$

Translational periodicity (block-circulant matrix)

Reduction to Block Diagonal Form

- In $3M \times 3M$ block form

$$K = \begin{bmatrix} \ddots & & & & & \\ & \left(K \begin{bmatrix} 0 & 0 \\ \alpha & \beta \end{bmatrix} \right) & & & & \\ & \left(K^T \begin{bmatrix} 0 & 1 \\ \alpha & \beta \end{bmatrix} \right) & & & & \\ & \left(K^T \begin{bmatrix} -1 & 1 \\ \alpha & \beta \end{bmatrix} \right) & & & & \\ & & \left(K \begin{bmatrix} 0 & 1 \\ \alpha & \beta \end{bmatrix} \right) & & & \\ & & \left(K \begin{bmatrix} 0 & 0 \\ \alpha & \beta \end{bmatrix} \right) & & & \\ & & \left(K^T \begin{bmatrix} 0 & 1 \\ \alpha & \beta \end{bmatrix} \right) & & & \\ & & & \left(K \begin{bmatrix} -1 & 1 \\ \alpha & \beta \end{bmatrix} \right) & & \\ & & & \left(K \begin{bmatrix} 0 & 1 \\ \alpha & \beta \end{bmatrix} \right) & & \\ & & & \left(K \begin{bmatrix} 0 & 0 \\ \alpha & \beta \end{bmatrix} \right) & & \\ & & & & \ddots & \end{bmatrix}$$

Translational periodicity (block-circulant matrix)
block-Fourier transform

$$\mathbb{K} = \begin{bmatrix} \ddots & & & & & \\ & \left(\mathbb{K} \begin{bmatrix} k-1 & \\ \alpha & \beta \end{bmatrix} \right) & & (0) & & (0) \\ & (0) & \left(\mathbb{K} \begin{bmatrix} k & \\ \alpha & \beta \end{bmatrix} \right) & & (0) & \\ & (0) & (0) & \left(\mathbb{K} \begin{bmatrix} k+1 & \\ \alpha & \beta \end{bmatrix} \right) & & \\ & & & & \ddots & \end{bmatrix}$$

- $3M \times 3M$ block diagonal form \implies Time complexity: $O\left([3M]^3 N\right)$.

- Crystal structure stability
 - Robust stability criterion — phonon spectra
 - Efficient numerical evaluation of phonon spectra
- Efficient equilibrium path following
 - Reduced set of equations based on symmetry
 - Pseudo-arc-length method
- Determine behavior near bifurcation points
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Following Equilibrium Paths

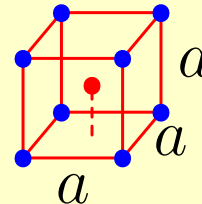
Solve a *reduced* set of equilibrium equations, e.g.,

• Cubic Phase

$$U_{11} = U_{22} = U_{33} = a,$$

$$U_{12} = U_{23} = U_{31} = 0.$$

Solve: $\frac{\partial \tilde{W}}{\partial a} = 0.$

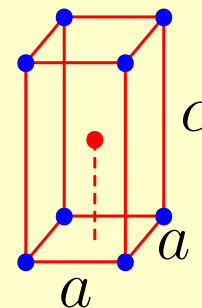


• Tetragonal Phase

$$U_{33} = c, \quad U_{11} = U_{22} = a,$$

$$U_{12} = U_{23} = U_{31} = 0.$$

Solve: $\frac{\partial \tilde{W}}{\partial a} = 0, \quad \frac{\partial \tilde{W}}{\partial c} = 0.$

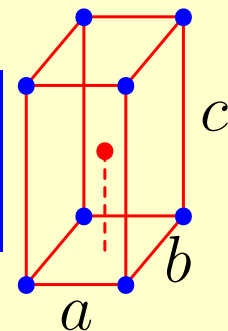


• Orthorhombic Phase

$$U_{11} = a, \quad U_{22} = b, \quad U_{33} = c,$$

$$U_{12} = U_{23} = U_{31} = 0.$$

Solve: $\frac{\partial \tilde{W}}{\partial a} = 0, \quad \frac{\partial \tilde{W}}{\partial b} = 0, \quad \frac{\partial \tilde{W}}{\partial c} = 0.$



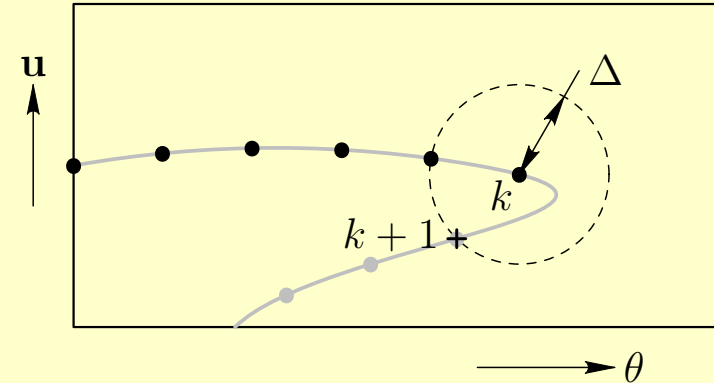
Advantages:

- Reduce computational effort
- **Eliminate singularities** near bifurcation points

Following Equilibrium Paths

Problem: following an equilibrium path around a turning point

- **Pseudo-arc-length method** (Riks method)
 - known solution $\mathbf{u}(\theta_k)$
 - find solution $\mathbf{u}(\theta_{k+1})$ a “distance” Δ away

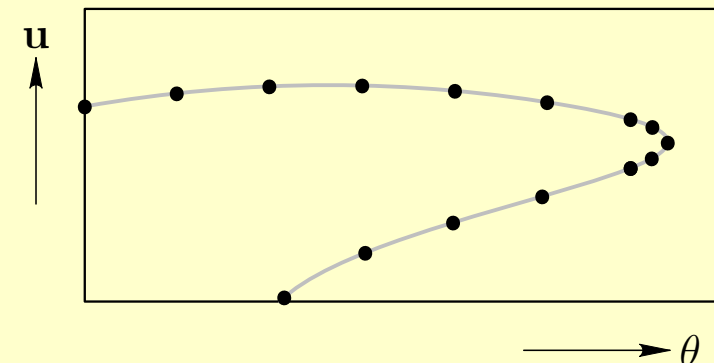


- Augment equilibrium equations with “distance” constraint

$$\frac{\partial \tilde{W}(\mathbf{u}(\theta_{k+1}); \theta_{k+1})}{\partial \mathbf{u}} = \mathbf{0}, \quad \|\mathbf{u}(\theta_{k+1}) - \mathbf{u}(\theta_k)\|^2 + (\theta_{k+1} - \theta_k)^2 = \Delta^2$$

- solve for θ_{k+1} and $\mathbf{u}(\theta_{k+1})$ simultaneously

- Also adaptively change Δ near turning points



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Asymptotic Bifurcation Analysis

- At a multiple bifurcation point, (\mathbf{u}_c, θ_c) :

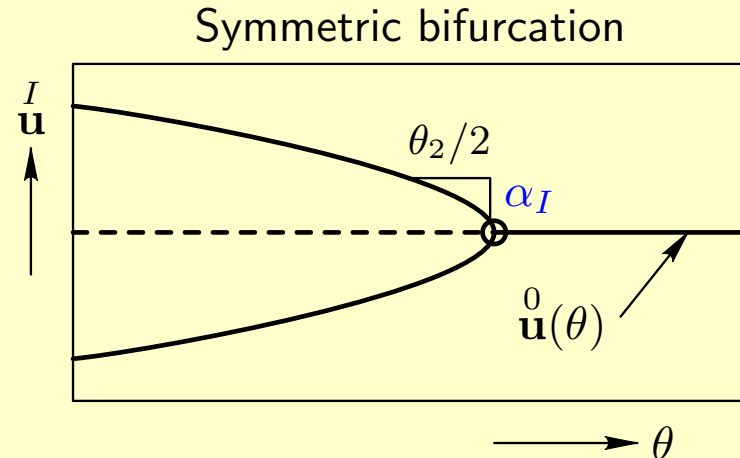
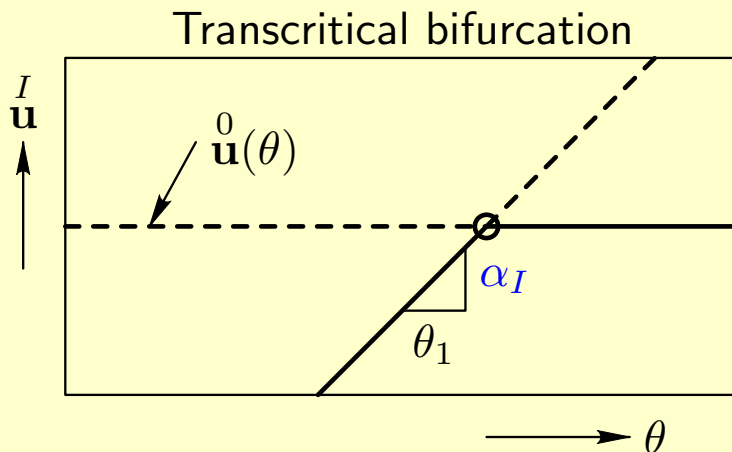
$\left. \frac{\partial^2 \tilde{W}}{\partial \mathbf{u}^2} \right|_c$ is singular with a null space of dimension $H \geq 2$.

- Following *Triantafyllidis & Peek (1992)*, (bifurcation amplitude parameter ξ)

$$\theta(\xi) = \theta_c + \theta_1 \xi + \theta_2 \frac{\xi^2}{2} + O(\xi^3),$$

$$\mathbf{u}(\xi) = \mathbf{u}^0(\theta(\xi)) + \left(\sum_{I=1}^H \alpha_I \mathbf{u}^I \right) \xi + \left(\sum_{I,J=1}^H \alpha_I \alpha_J \mathbf{v}^{IJ} \right) \frac{\xi^2}{2} + O(\xi^3),$$

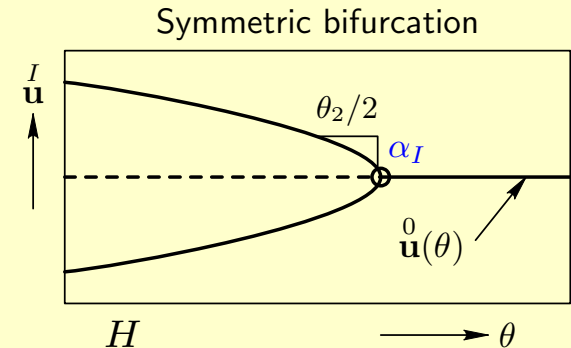
where $\{\mathbf{u}^1, \dots, \mathbf{u}^H\}$ is an O.N. basis for the null space of $\left. \frac{\partial^2 \tilde{W}}{\partial \mathbf{u}^2} \right|_c$.



Substitute into equilibrium equations $\frac{\partial \tilde{W}}{\partial \mathbf{u}} = \mathbf{0}$, expand, and collect L.O.T.

- For *symmetric* bifurcation: $\left(\frac{\partial^3 \tilde{W}}{\partial \mathbf{u}^3} \Big|_c \begin{matrix} I & J & K \\ \mathbf{u} & \mathbf{u} & \mathbf{u} \end{matrix} \right) \equiv 0$

$$\sum_{J,K,L=1}^H \alpha_J \alpha_K \alpha_L \mathcal{E}_{IJKL} + 3\theta_2 \sum_{J=1}^H \alpha_J \mathcal{E}_{IJ\theta} = 0,$$



$$\sum_{I=1}^H (\alpha_I)^2 = 1.$$

$$\mathcal{E}_{IJ\theta} \equiv \left(\frac{d}{d\theta} \left(\frac{\partial^2 \tilde{W}(\mathbf{u}^0(\theta); \theta)}{\partial \mathbf{u} \partial \mathbf{u}} \right) \right) \Big|_c \begin{matrix} I & J \\ \mathbf{u} & \mathbf{u} \end{matrix},$$

$$\mathcal{E}_{IJKL} \equiv \frac{\partial^4 \tilde{W}}{\partial \mathbf{u}^4} \Big|_c \begin{matrix} J & K & L & I \\ \mathbf{u} & \mathbf{u} & \mathbf{u} & \mathbf{u} \end{matrix} + \frac{\partial^3 \tilde{W}}{\partial \mathbf{u}^3} \Big|_c \left(\begin{matrix} J & K & L \\ \mathbf{v} & \mathbf{u} \end{matrix} + \begin{matrix} K & L & J \\ \mathbf{v} & \mathbf{u} \end{matrix} + \begin{matrix} L & J & K \\ \mathbf{v} & \mathbf{u} \end{matrix} \right) \begin{matrix} I \\ \mathbf{u} \end{matrix},$$

$$\frac{\partial^2 \tilde{W}}{\partial \mathbf{u}^2} \Big|_c \begin{matrix} I & J \\ \mathbf{v} & \mathbf{u} \end{matrix} = - \frac{\partial^3 \tilde{W}}{\partial \mathbf{u}^3} \Big|_c \begin{matrix} I & J \\ \mathbf{u} & \mathbf{u} \end{matrix}$$

- Fredholm alternative guarantees a unique $\begin{matrix} I & J \\ \mathbf{v} & \mathbf{u} \end{matrix}$

Need $\frac{IJ}{\mathbf{v}}$:
$$\left. \frac{\partial^2 \tilde{W}}{\partial \mathbf{u}^2} \right|_c \frac{IJ}{\mathbf{v}} = - \left. \frac{\partial^3 \tilde{W}}{\partial \mathbf{u}^3} \right|_c \frac{IJ}{\mathbf{u}\mathbf{u}}$$

- Generate an O.N. basis for \mathbb{R}^n by diagonalizing $\left. \frac{\partial^2 \tilde{W}}{\partial \mathbf{u}^2} \right|_c \in \mathbb{R}^n \times \mathbb{R}^n$

$$\mathcal{N} = \text{Span} \left\{ \mathbf{1}, \dots, \mathbf{u}^H \right\}, \quad \mathcal{N}^\perp = \text{Span} \left\{ \mathbf{v}, \dots, \mathbf{v}^{n-H} \right\}$$

- Projection operator $[Q_{IJ}] = [\mathbf{v}_j^I] : \mathbb{R}^n \mapsto \mathcal{N}^\perp$

$$\underbrace{Q \left. \frac{\partial^2 \tilde{W}}{\partial \mathbf{u}^2} \right|_c Q^T}_{\text{non-singular}} Q \frac{IJ}{\mathbf{v}} = -Q \left. \frac{\partial^3 \tilde{W}}{\partial \mathbf{u}^3} \right|_c \frac{IJ}{\mathbf{u}\mathbf{u}}$$

- Solving gives

$$\frac{IJ}{\mathbf{v}} = \underbrace{-Q^T}_{n \times (n-H)} \underbrace{\left[Q \left. \frac{\partial^2 \tilde{W}}{\partial \mathbf{u}^2} \right|_c Q^T \right]^{-1}}_{(n-H) \times (n-H)} \underbrace{Q \left. \frac{\partial^3 \tilde{W}}{\partial \mathbf{u}^3} \right|_c \frac{IJ}{\mathbf{u}\mathbf{u}}}_{(n-H) \times 1}$$

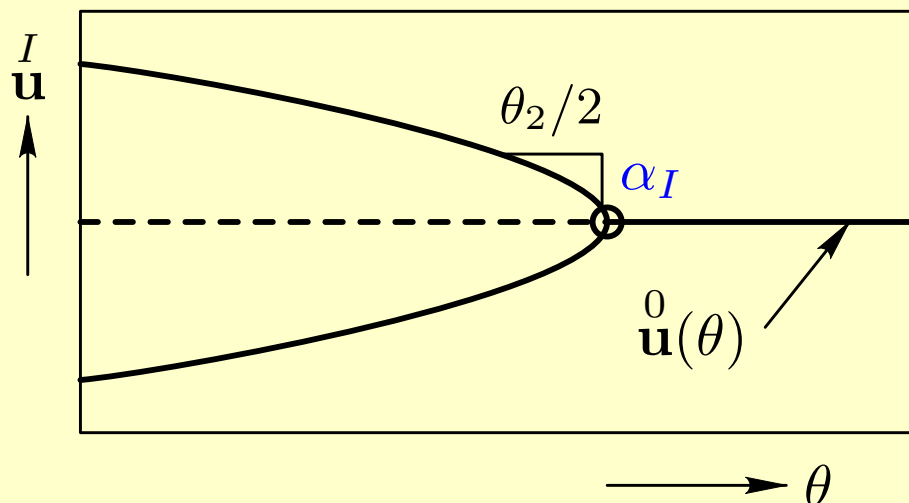
$n \times 1$

- $\mathcal{E}_{IJ\theta}$, \mathbf{v}^{IJ} , and \mathcal{E}_{IJKL} are obtained numerically
- All bifurcating equilibrium paths are found by solving

$$\sum_{J,K,L=1}^H \alpha_J \alpha_K \alpha_L \mathcal{E}_{IJKL} + 3\theta_2 \sum_{J=1}^H \alpha_J \mathcal{E}_{IJ\theta} = 0, \quad \sum_{I=1}^H (\alpha_I)^2 = 1.$$

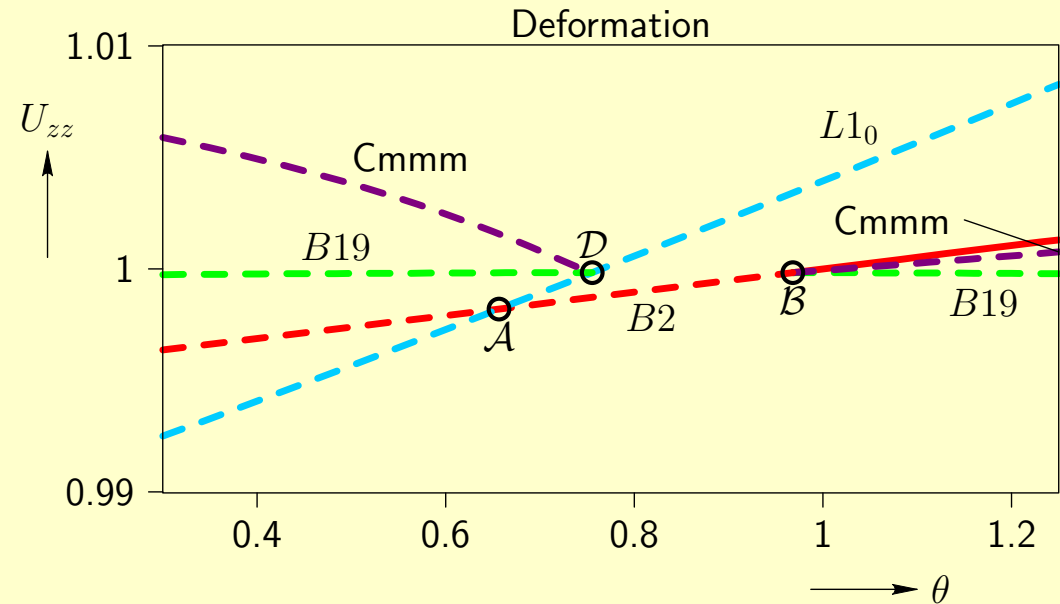
- In general there are $(3^H - 1)/2$ pairs of solutions (α_I, θ_2) and $(-\alpha_I, \theta_2)$
- Each pair of solutions corresponds to one symmetric equilibrium path

Symmetric bifurcation



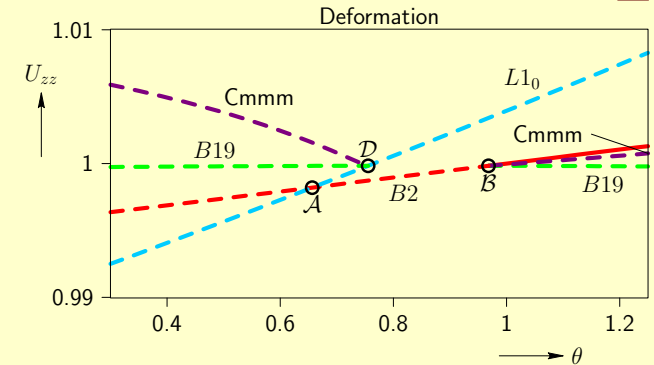
Example of Degree Two Bifurcation

- Bifurcation of degree two ($H = 2$) at \mathcal{B}



Example of Degree Two Bifurcation

- Bifurcation of degree two ($H = 2$) at \mathcal{B}



- Basis for \mathcal{N} (translation of certain crystal planes)

$$\mathbf{u}^1 = \begin{bmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0.6931 & 0 & 0 & 0.7203 & 0 & 0 & 0.0271 & 0 \end{bmatrix},$$

$$\mathbf{u}^2 = \begin{bmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0.0271 & 0 & 0 & 0 & 0.7203 & 0 & 0 & 0.6931 & 0 & 0 \end{bmatrix}$$

- Resulting bifurcation equations ($\theta_1 = 0$)

$$-4849.2(\alpha_1)^3 - 19421(\alpha_1(\alpha_2)^2) + 3\theta_2(0.02425\alpha_1) = 0,$$

$$-19421((\alpha_1)^2\alpha_2) - 4849.2(\alpha_2)^3 + 3\theta_2(0.02425\alpha_2) = 0.$$

- Solutions (four pairs)

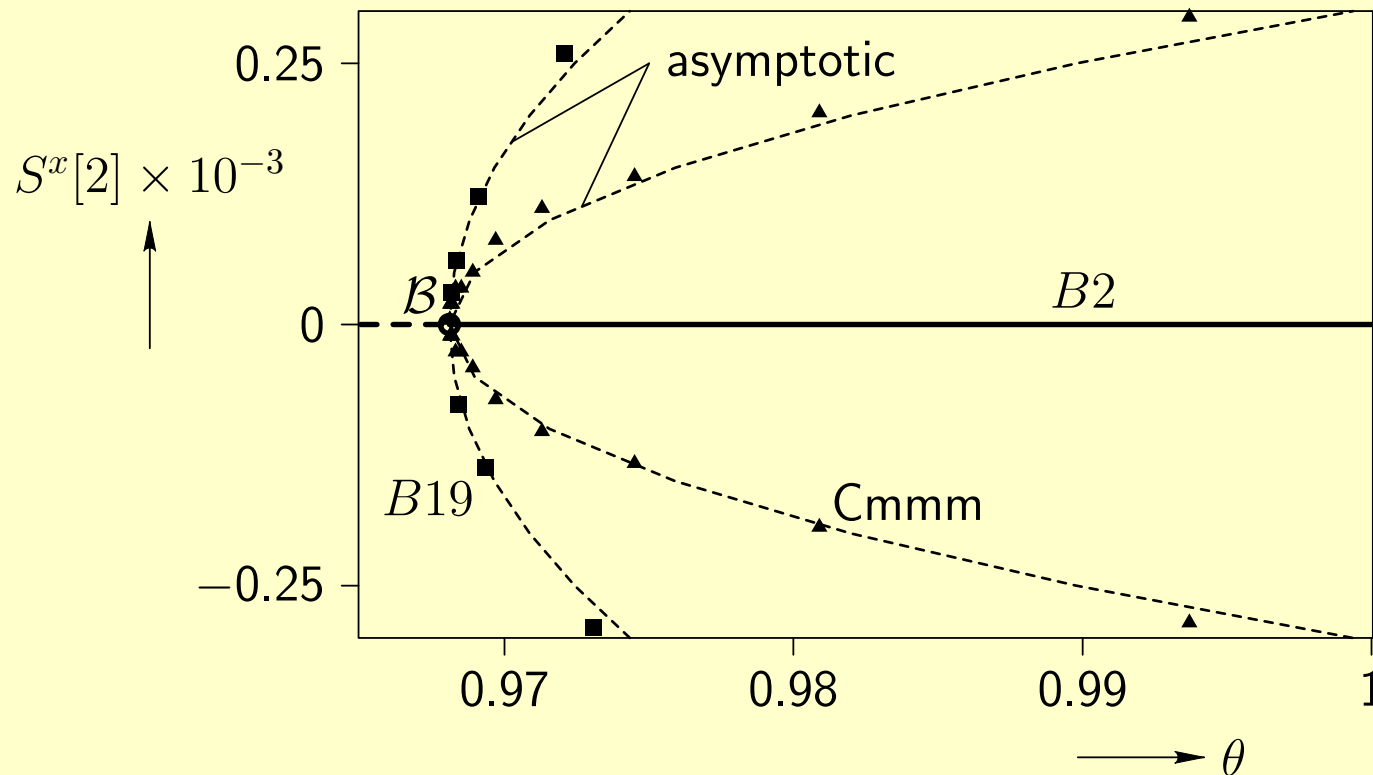
$$B19(1) : \quad \alpha_1 = 1, \alpha_2 = 0, \quad \theta_2 = 66654,$$

$$B19(2) : \quad \alpha_1 = 0, \alpha_2 = 1, \quad \theta_2 = 66654,$$

$$Cmmm(1) : \quad \alpha_1 = 1, \alpha_2 = 1, \quad \theta_2 = 333607,$$

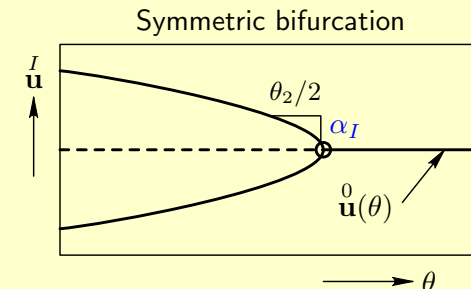
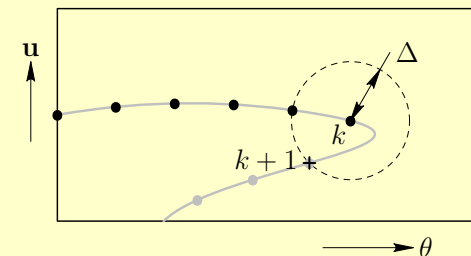
$$Cmmm(2) : \quad \alpha_1 = 1, \alpha_2 = -1, \quad \theta_2 = 333607.$$

- Compare numerical and asymptotic results



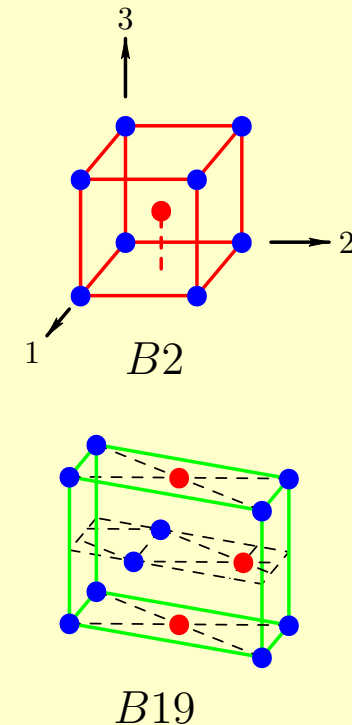
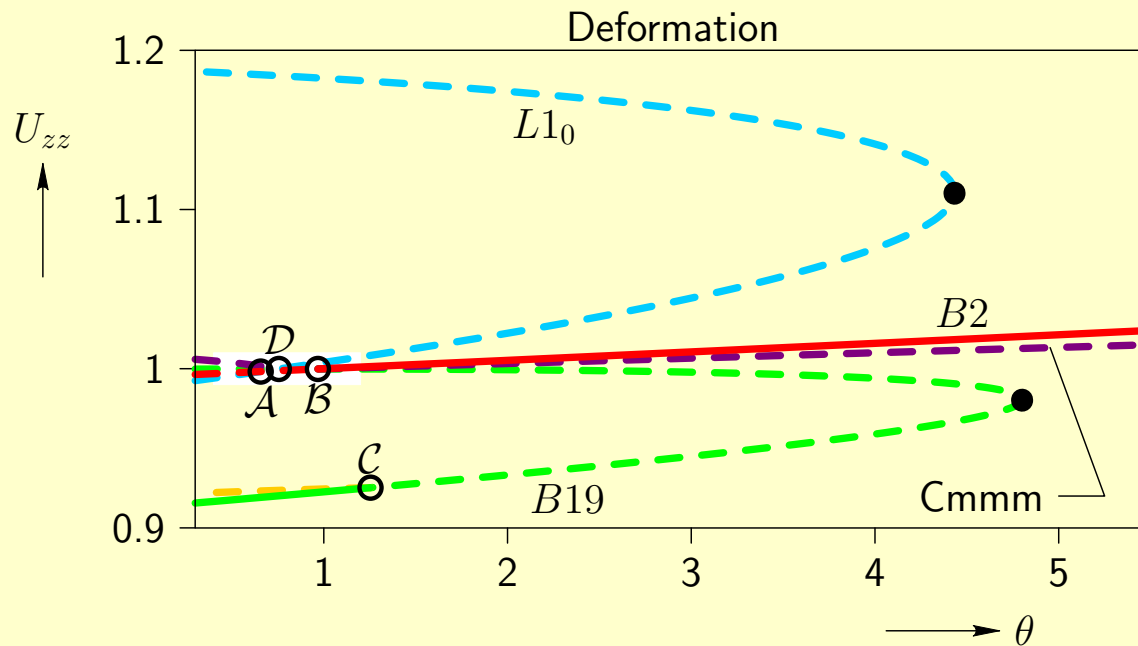
Numerical techniques for bifurcation investigation of atomistic material models

- Phonon spectra — important measure of crystal stability
 - Block-Fourier transform allows efficient phonon spectra computation
 - Time complexity $O((3M)^3 N)$: linear in number of unit cells
- Efficient methods for following equilibrium paths
 - Reduce the number of equations by invoking symmetry
 - Pseudo-arc-length method
- Analyze behavior near bifurcation points
 - Numerically assisted asymptotic bifurcation investigation



Summary & Conclusions

- Used these computational techniques to study a new atomistic model
 - Temperature-dependent atomic-potentials
 - Cauchy-Born kinematics — uniform deformation & internal shifts



- Identified a **hysteretic proper Martensitic transformation**
 - Cubic austenite phase ($B2$ CsCl-type crystal)
 - Orthorhombic martensite phase ($B19$ crystal structure)
 - These structures are experimentally observed in SMA's such as AuCd, CuAlNi, and NiTiCu.