Less is More: Efficient Numerics for Model Physics Problems in Simple Geometries

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Features of Model Problems in Physics & Engineering

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- Focus on dominant physics
- Geometry is *simple*

y

X



Pseudospectral Methods

- Advantages
 - Very computationally efficient:

error ~ $O(p^N)$, $p < 1 \longrightarrow$



- Implementation effort similar to FD
- Disadvantages
 - Difficult to handle irregular geometries
 - Effort needed to handle non-smooth solutions

Two Basic Ideas

- Choose grid points appropriately based on computational domain



• Use **ALL** grid points to compute derivatives

Fully dense differentiation matrices ... for SMALL $N\,!$

Computational Domains & Grids

• Periodic

$$x_k = \frac{2k - N}{N}$$

- Finite, non-periodic $x_k = \cos(k\pi/N)$
- Infinite

$$y_k = \frac{Lx_k}{\sqrt{1 - x_k^2}}$$

• Semi-infinite

$$y_k = L\left(\frac{1-x_k}{1+x_k}\right)$$



From FD to PS

- I. Uniform grid **—** pseudospectral grid
- Low-order discrete derivative
 pseudospectral differentation matrix
- 3. Use standard numerical methods (e.g. Newton iteration, discrete time-stepping)

"Inside every low-order program is a high order algorithm waiting to burst free!" – J.P. Boyd

Direct Matrix Method for Computing Exact Jacobians

- Write discretized differential equations **explicitly** in terms of differentation matrices
- Examples:

$$\frac{\partial^2 \phi}{\partial x^2} = \rho \implies D^2 \phi = \rho$$

$$(\nabla\phi) \cdot (\nabla\psi) + \psi\nabla^2\phi = 0 \longrightarrow$$
$$(D_x\phi) \cdot (D_x\psi) + (D_y\psi) \cdot (D_y\psi) + \psi \cdot (L\phi) = 0$$

Direct Matrix Method for Computing Exact Jacobians

• Simple differentiation rules to compute **exact** Jacobian

$$\begin{aligned} \frac{\partial}{\partial u}(A * u) &= A\\ \frac{\partial f(u)}{\partial u} &= \text{diag}\left[f'(u)\right]\\ \frac{\partial}{\partial u}\left[A * f(u)\right] &= A * \text{diag}\left[f'(u)\right]\\ \frac{\partial}{\partial u}\left[f(u) \cdot g(u)\right] &= \text{diag}\left[g(u)\right] * \frac{\partial f}{\partial u} + \text{diag}\left[f(u)\right] * \frac{\partial g}{\partial u}\end{aligned}$$

Less tedious and less error-prone Jacobian calculation!

Direct Matrix Method for Computing Exact Jacobians

• Example

$$F(u) = u'' + \sin(2u)u'$$

$$F(u) = D^2 * u + \sin(2u) * (D * u)$$

 $\frac{\partial F}{\partial u} =$

 $D^{2} + \operatorname{diag}[\sin(2u)] * D + 2\operatorname{diag}[D * u] * \operatorname{diag}[\cos(2u)]$

Making the Most of MATLAB

- Advantages
 - Many useful built-in functions
 - Easy visualization
- Disadvantage: considered slow
 - Solution: mixed-language programming (i.e. MEX-files)

High-efficiency programming + High-performance computing

Electrochemical Thin-Films

• Integro-differential equation

$$\epsilon^{2} \left[\frac{d^{2}E}{dx^{2}} - \frac{1}{2}E^{3} \right] - (c_{0} + 2jx)E = 2j$$

$$c_0 = (1-j) + \epsilon^2 \left[E(1) - E(0) - \frac{1}{2} \int_0^1 E^2 dx \right]$$

• Nonlinear reaction boundary conditions

$$k_c c_+ e^{-\alpha_c \delta \epsilon E(0)} - j_r e^{\alpha_a \delta \epsilon E(0)} = j$$
$$-k_c c_+ e^{\alpha_c \delta \epsilon E(1)} + j_r e^{-\alpha_a \delta \epsilon E(1)} = j$$

Electrochemical Thin-Films



Charging of Metal Colloid Sphere



- Governing equations $0 = \nabla^2 c \qquad 0 = \nabla \cdot (c \nabla \phi)$
- Boundary conditions $F_s^{\pm} = -\frac{\epsilon \left(e^{\pm \zeta/2} - 1\right)}{\sqrt{c}} \left(\nabla_s c \pm 2c \nabla_s \phi\right) \qquad \zeta + 2\delta(\zeta/2) = -\phi$

Charging of Metal Colloid Sphere







Conformal Mapping for Nernst-Planck Equations

• Steady, electroneutral Nernst-Planck equations are conformally invariant (Bazant 2004)



High-Performance **Desktop** Scientific Computing

- Spectral methods yield *high* accuracy at *low* cost
- **Exact** Jacobians via Direct Matrix Method
- MATLAB + C/Fortran = powerful tool for moderate-sized problems

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THE END

EXTRA SLIDES

Desktop "Super"-Computers

- Fast CPUs (> 3 GHz)
- Large memories (2 4 GB)
- Large on-chip caches (1024 K)

Modern serial computers are extraordinarily powerful!